

# Introduction to ATLAS

- part 1: ATLAS Detector (and LHC)
- part 2: Physics programme in ATLAS
- **part 3: Event Reconstruction and Physics Performance**
- part 4: Physicists' tools analyses in ATLAS



*Wolfgang Liebig*

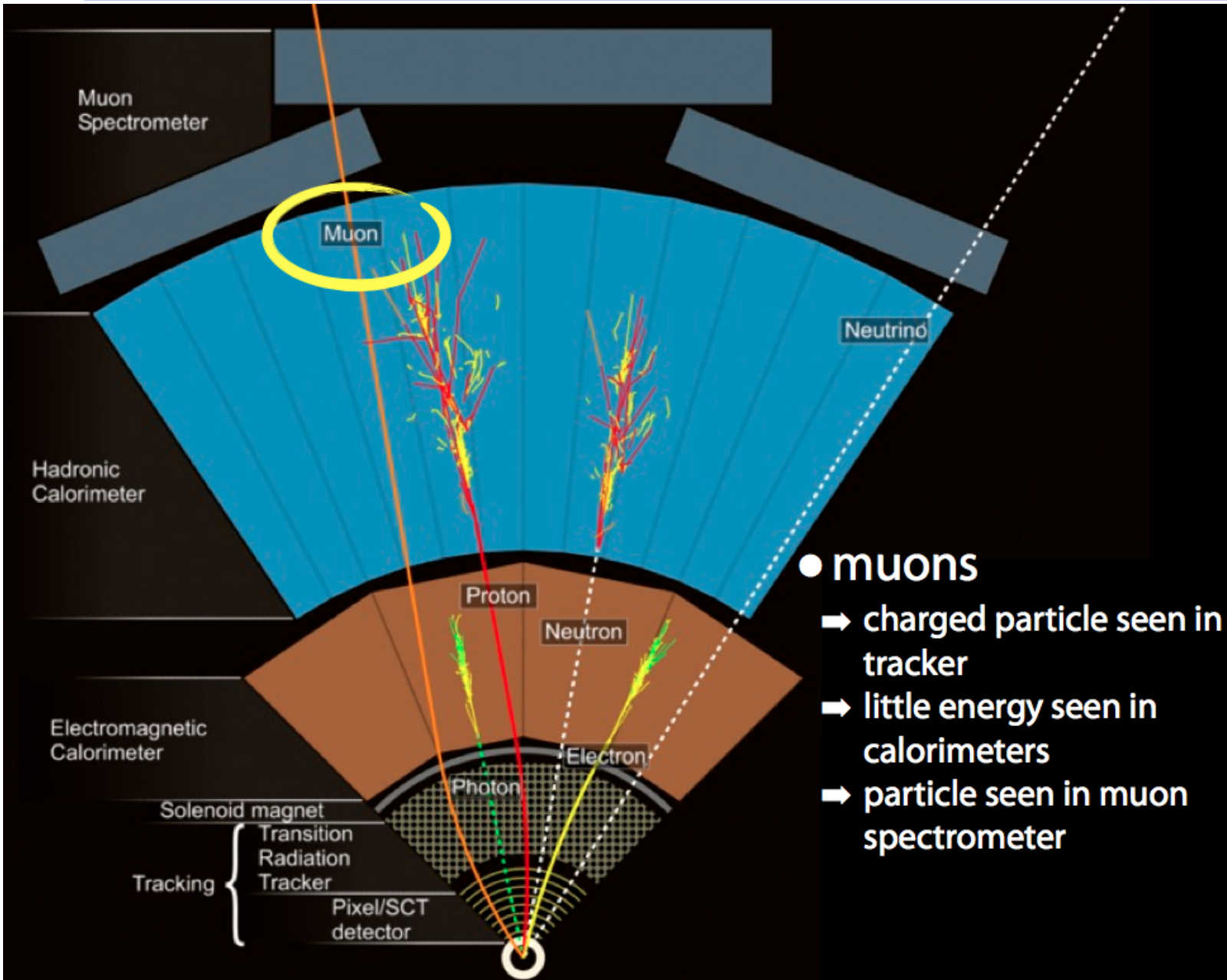


# Part 3: Reconstruction & Performance

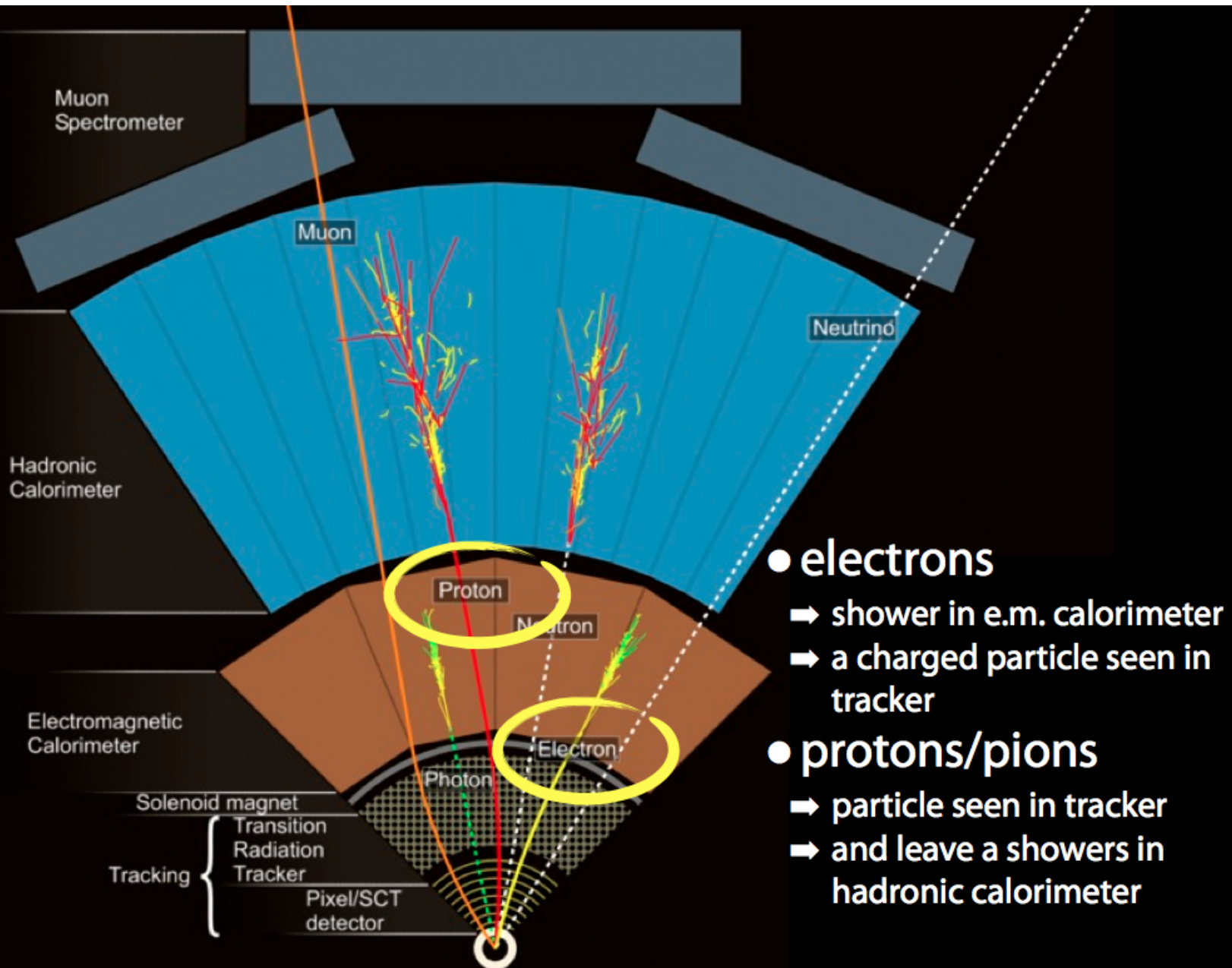
- Physics objects – overview
- Tracking
  - track representation, fit, performance
- muon performance
  - identification, performance, alignment
- vertex reconstruction
- b-tagging
- electron and photon reconstruction and performance
- jet reconstruction
- missing transverse energy
- tau lepton reconstruction



# Reconstruction in a Nutshell

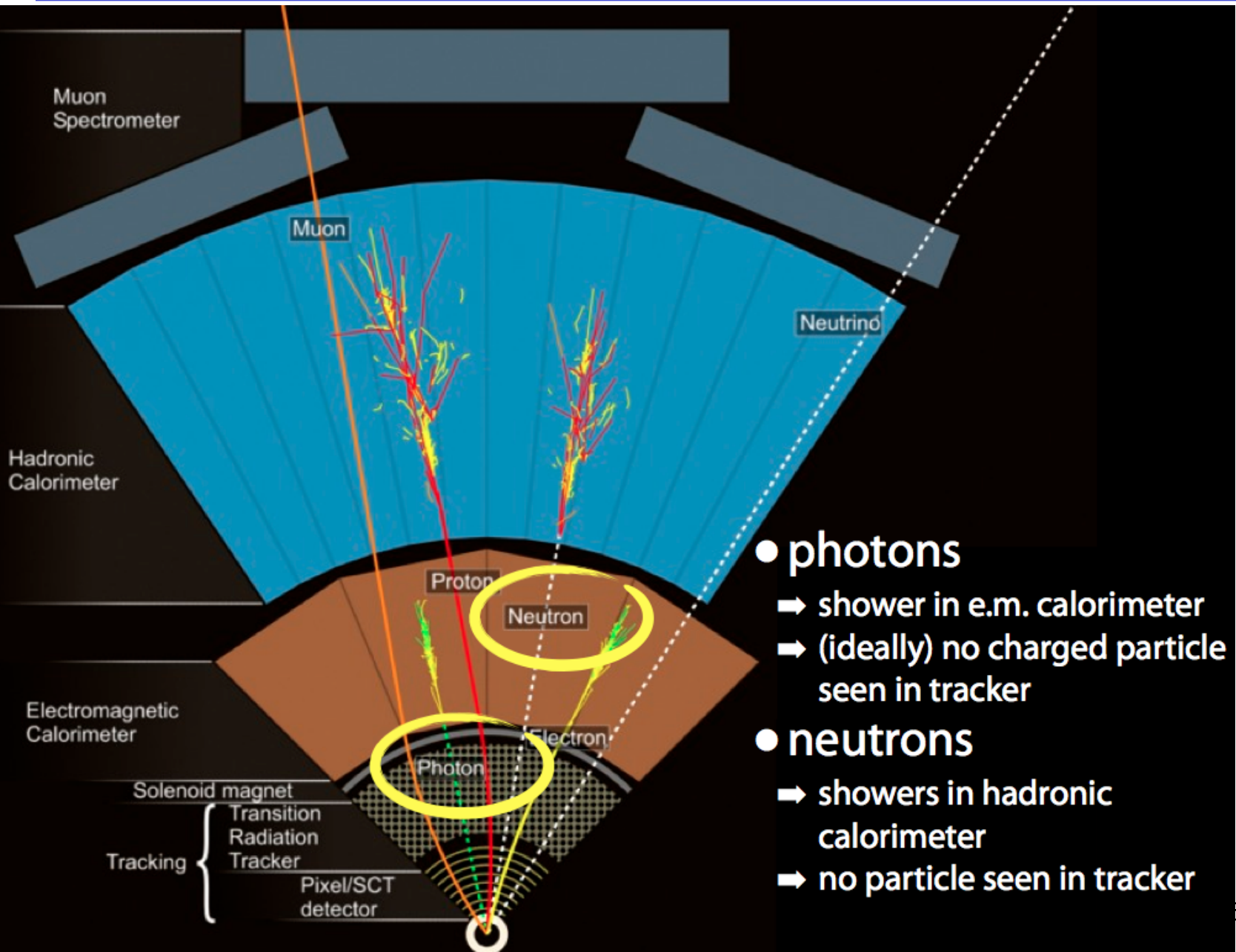


# Electrons and Hadrons

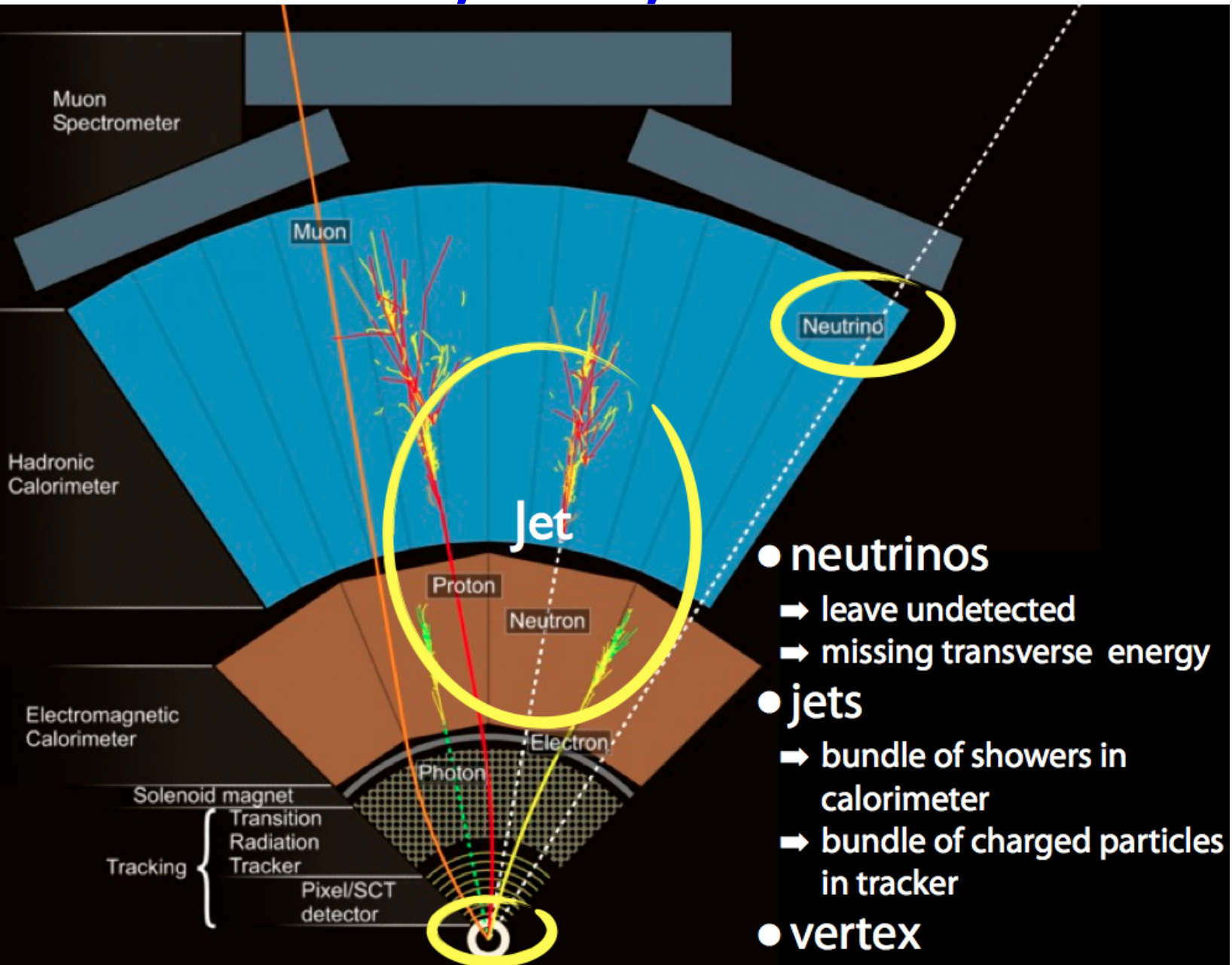




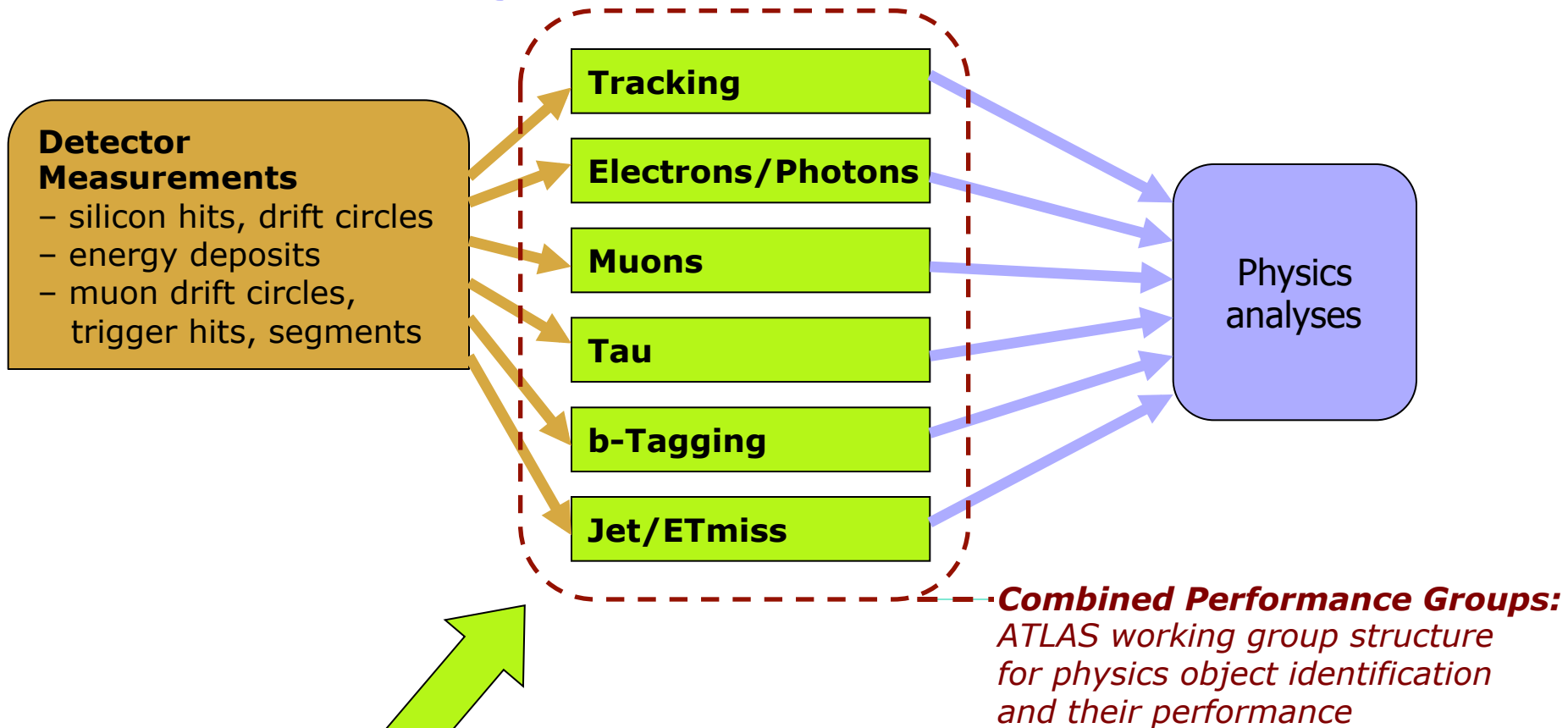
# Photons and Neutrons



# Neutrinos, Jets, Vertex



# Physics Objects in ATLAS



- Performance
  - Efficiency of object identification
  - Purity or fraction of false positive
  - Energy or momentum resolution
  - recommendations for physics analyses
  - combined data quality for physics object

- Tracking is input to all other
- Trigger
  - similar but separate group structure
  - trigger efficiency is main performance number



# Performance: Basic Concepts

- **Efficiency** of identification
  - $N_{\text{true}}(\text{identified})/N_{\text{true}}(\text{all})$
  - multiplicative for events with several physics objects
- efficiency determination
  - simulation: MC truth tells  $N_{\text{true}}(\text{all})$
  - data: methods using known physics and/or detector redundancy
- **Trigger Efficiency**
  - of event passing trigger (  $\rightarrow$  stream)
  - of offline physics object (e.g. muon) being the one that fired trigger
- trigger eff. determination
  - simulation: combine L1–L2–EF
  - data: use offline object as 'truth'
- **Isolation efficiency**
  - object identified with additional isolation requirement
- **Fake Rate** of identification
  - $N(\text{misidentified})/N(\text{identified})$
  - simulation: MC truth tells misidentification
  - data: identification counts on sample known to be depleted of physics object
- **Resolution** of track parameters
  - average difference between true and reconstructed parameters
  - most common: momentum and impact parameter resolutions
  - simulation: MC truth for parameters
  - data: knowledge about physics process, detector redundancy etc.
- ideally resolution should be reflected by error on parameters



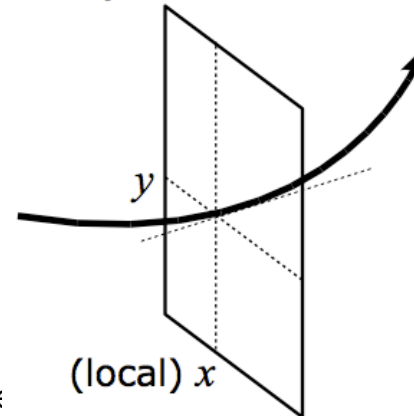
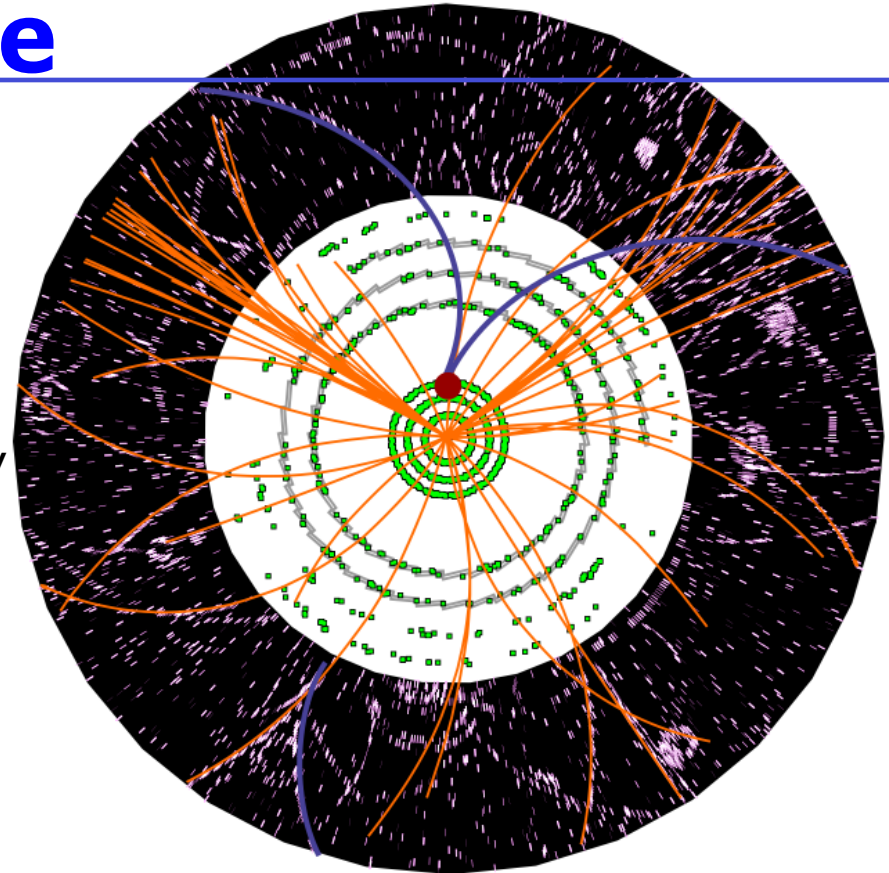
# Tracking, Muons, B-tagging...



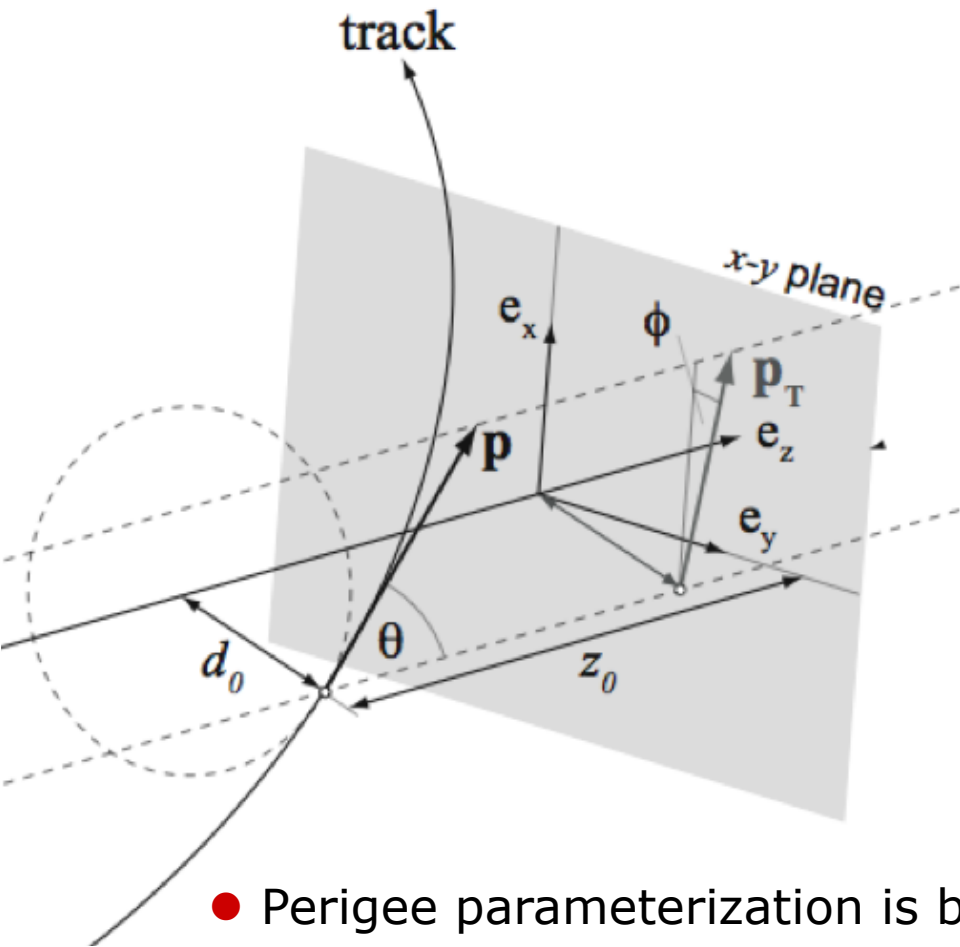


# Tracking Software

- Charged particles leave a “cloud of hits” in the detector
  - further obscured by hits from noise, interactions with detector material, low energy curling tracks.
- Tracking software needs to identify particle trajectories, reconstruct their kinematic parameters
- Track model parameterizes trajectory with 5 parameters
  - stable particle moving in stationary B-field in vacuum described by 6 par
  - initial position along trajectory is free
- Local Parameters of track model
  - at an intersected reference surface:
    - 2 local coordinates
    - 2 angles
    - curvature  $q/p$
  - and their 5x5 covariance matrix



# Track Parameters at Collider Detectors



## Perigee parameterization:

- $d_0$  signed distance of closest approach to  $z$  axis
- $z_0$   $z$  of closest approach
- $\phi_0$  azimuthal angle at cl. app.
- $\theta$  polar angle of track
- $q/p$  charge-signed curvature

- Perigee parameterization is basis for
  - expressing track parameters at production vertex for instance Lorentz vectors in physics analysis
  - vertex finding algorithms
  - b-jet tagging



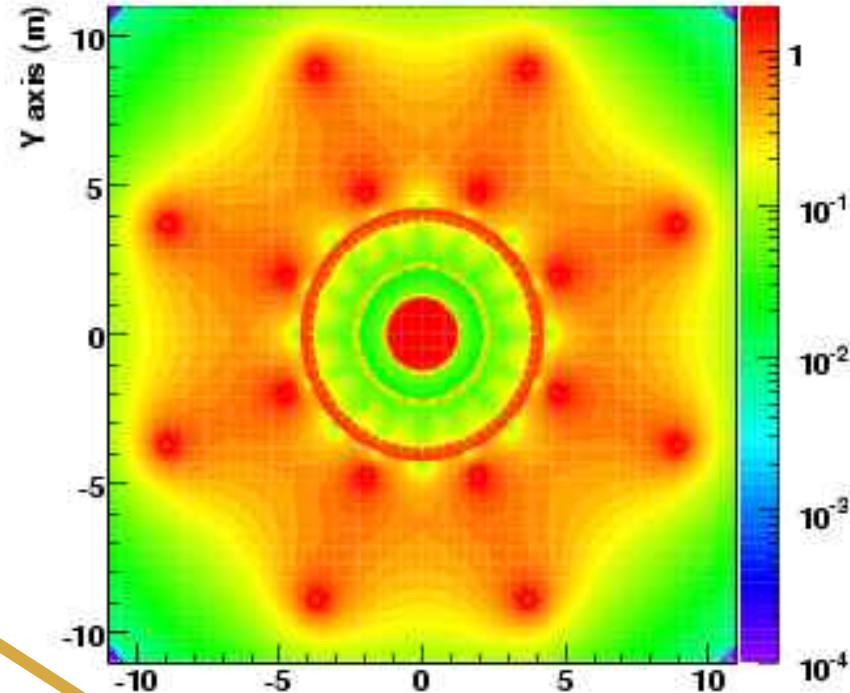
# Track Propagation: Fields

- **Equation of Motion** of particle

$$\frac{d^2 \mathbf{r}}{ds^2} = \frac{q}{p} \left( \frac{d\mathbf{r}}{ds} \times \mathbf{B}(\mathbf{r}) \right)$$

- helix approximation not sufficient:
  - risk  $\sim 1\%$  momentum bias (CMS?)
  - ATLAS InDet longer than solenoid
  - toroids produce inhomogeneous field
- $\mathbf{B}(\mathbf{r})$  inhomogeneous: diff. equation can only be solved numerically
- Runge-Kutta-Nystrom methods
  - divide integration interval in steps
  - each step becomes initial-value problem
  - solve equation for each step individually
- form the detailed track model
  - in ATLAS called **Propagator**
  - model of interactions in detector added separately

z=-20cm, phi=2pi



**common tracking software**  
designed to work in both  
Inner Detector  
and Muon Spectrometer

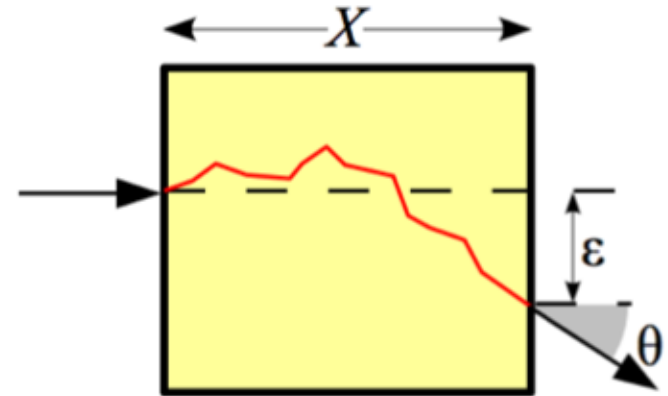


# Multiple Scattering

Besides field effects, track propagation also affected by material:

- energy loss (discussed in part 1 of lecture)
- **multiple scattering**

- charged particle **deflected** when passing through matter
- random deflection is result of many small-angle Coulomb scatterings on the nuclei
- Gaussian distribution for central 98 % given by **Highland formula**



$$\sigma(\theta) = \frac{13.6 \text{ MeV}}{\beta c p} z \sqrt{x/X_0} (1 + 0.038 \ln(x/X_0))$$

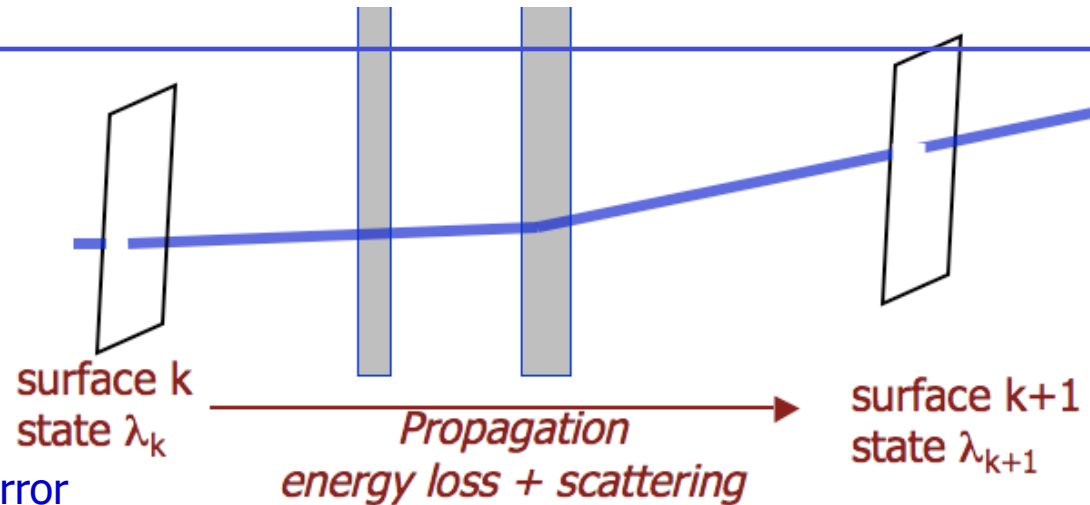


- expect  $E(\epsilon) = 0$ ,  $E(\theta) = 0$ .  $\sigma(\theta)$  is proportional to  $1/p$
- $x/X_0$ : thickness of material in **fraction of radiation length**

# Track Fitting

- Measurement constraints

- $m_k = h_k(\lambda) + \gamma_k$
- $\lambda$ : track parameter vector  
(propagation to be added)
- $h_k$ : functional dependence of measurement on track parameters (meas't model)
- $\gamma_k$ : noise term, variation within error



- Now need an estimator for  $\lambda$

- could e.g. use MINUIT (max. likelihood)  
but that is not the case in tracking

- A linear model is applied

- $h_k(\lambda) \simeq h_k^0 + H_k \cdot \lambda$
- $H_k = \frac{dm_k}{d\lambda}$  : Jacobian, typically a rotation or projection into measurement plane

- Measurements are Gaussian distributed

- least squares estimator is best unbiased estimator

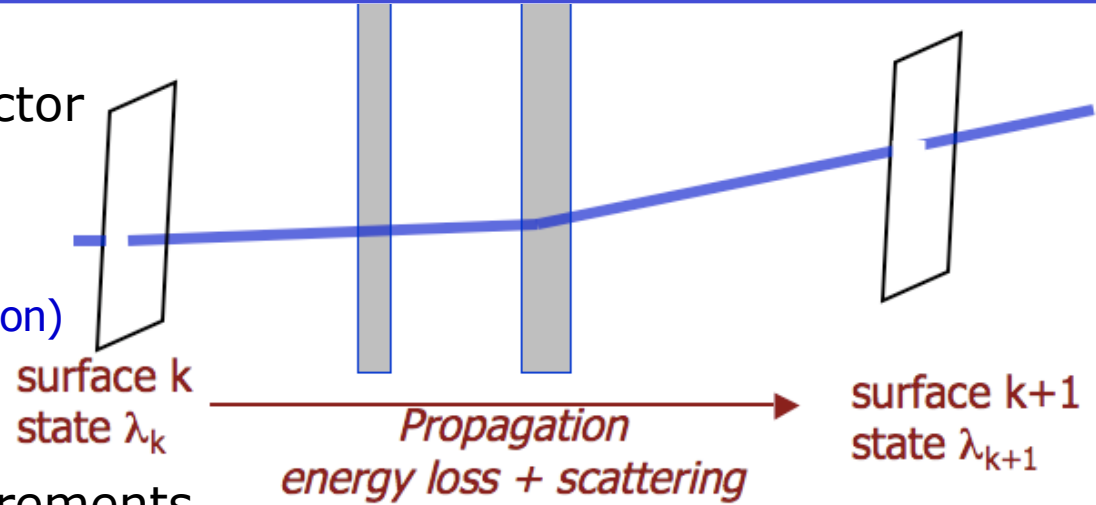
$$\chi^2 = \sum_i^{N_{\text{hits}}} \left( \frac{m_i - h_i(\lambda, \theta^{\text{scat}})}{\sigma_i} \right)^2 + \sum_j^{N_{\text{planes}}} \left( \frac{E(\theta^{\text{scat}}) - \theta_j^{\text{scat}}}{\sigma^{\text{scat}}} \right)^2$$





# Track Fitting in ATLAS

- Propagators transport  $\lambda$  vector
  - $\lambda_k^{k-1} = f_k^{k-1}(\lambda_{k-1})$
  - simplified geometry used (simpler+ faster than full simulation)
  - algorithm “AtlasExtrapolator”
- “Global  $\chi^2$  fitters” solve lin. estimator for all measurements
  - needs inversion of large matrix dim=5+2N
  - fit follows trajectory closely, useful for large distances in ATLAS
  - mostly used in ATLAS
- Kalman filters in track fitting
  - steps through hits and updates parameters
  - progressive way of performing LSE, mathematically equivalent
  - fast: series of dim.5 matrix inversions
  - extended to estimate electron trajectories (bremsstrahlung, in use in ATLAS)

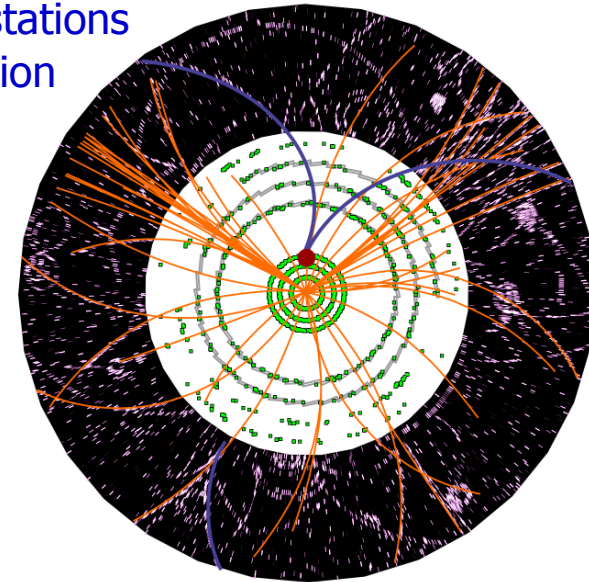
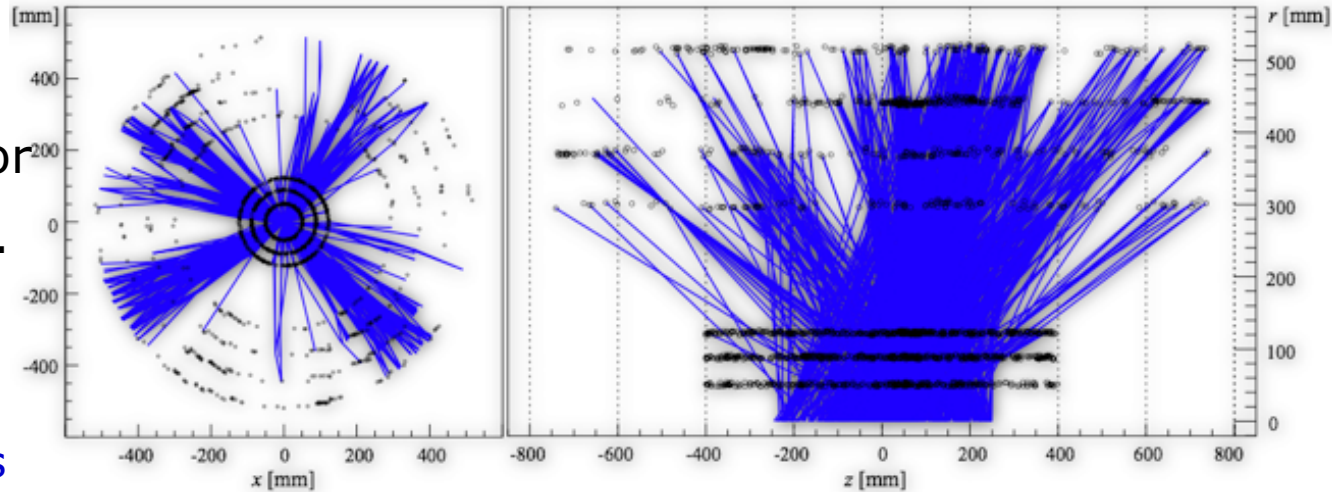


- Robust estimators
  - define ‘pull’ =  $|m_k - \hat{h}_k(\lambda)| / \sigma_k$
  - typically reject hit as an outlier if pull > 3.5
  - re-fit with outliers rejected if  $\text{prob}(\chi^2) < 10^{-5}$
  - avoids bias or degraded resolution



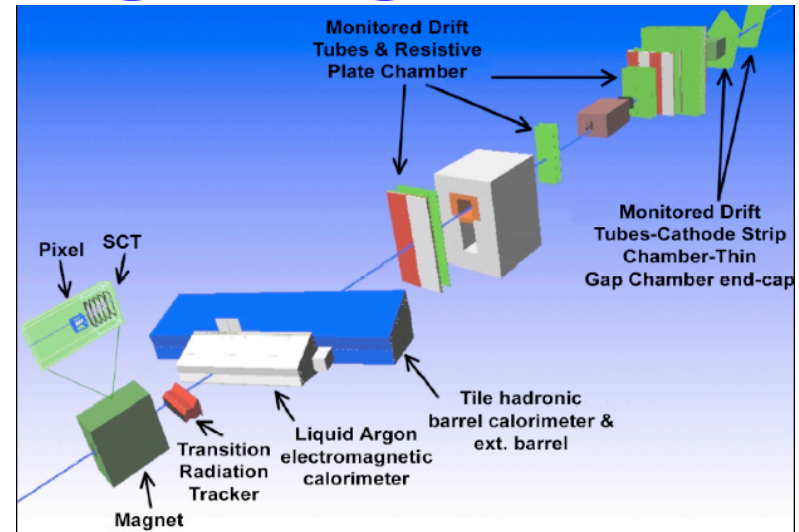
# Track Finding

- Choice of track finding algorithm depends on detector
  - robustness against combinatoric problems and detector ambiguities needed
  - algorithms use Hough-transforms or look-up tables
- Seed finding in e.g. silicon, TRT, muon chambers
  - seeds or segments extended by combinatorial track following
  - associates hits in adjacent layers or muon segments in other stations
  - upon ambiguities branch seed following and evaluate best option
  - often fast versions of track fit employed in track following
- combination of inside-out (for prompt tracks) and outside-in seeding (e.g.  $\gamma$  conversions)
- final track fit with precise material corrections and hit recalibration

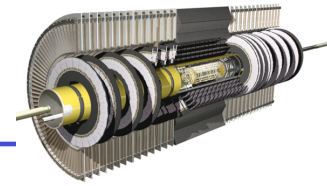


# ATLAS Commissioning Programme

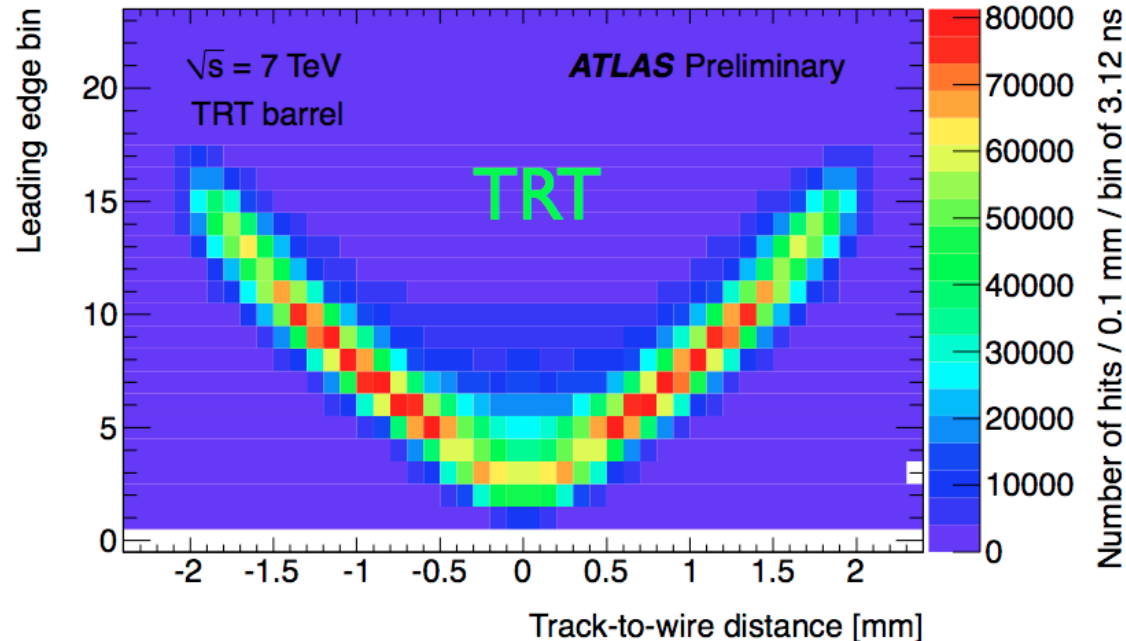
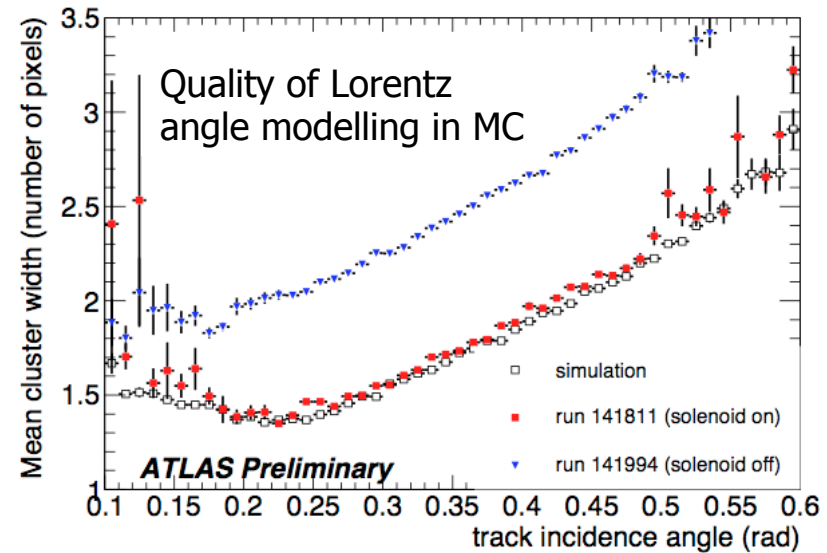
- 1995-2004 test beams  
2004 combined test beam →
  - software integration
  - first performance measurements for single particle detection
- 2006 cosmic rays (Inner Detector)  
2008/9 cosmic rays (whole ATLAS as installed in P1)
  - first performance measurements on real detector (tracks+muons)
- 2009 single beam events and 900 GeV collisions
  - correct some obvious performance mis-modelling (dead modules, missing large structures in detector geometry)
- 2010/11 p-p collisions
  - high statistics of tracks “illuminate” even remote MS chambers
  - high statistics of calibration objects, such as  $J/\psi$  and Z decays to  $e, \mu, \tau$
  - methods often tricky and similarly involved as a physics analysis



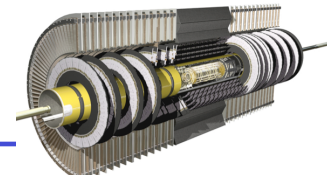
# Inner Detector Calibration



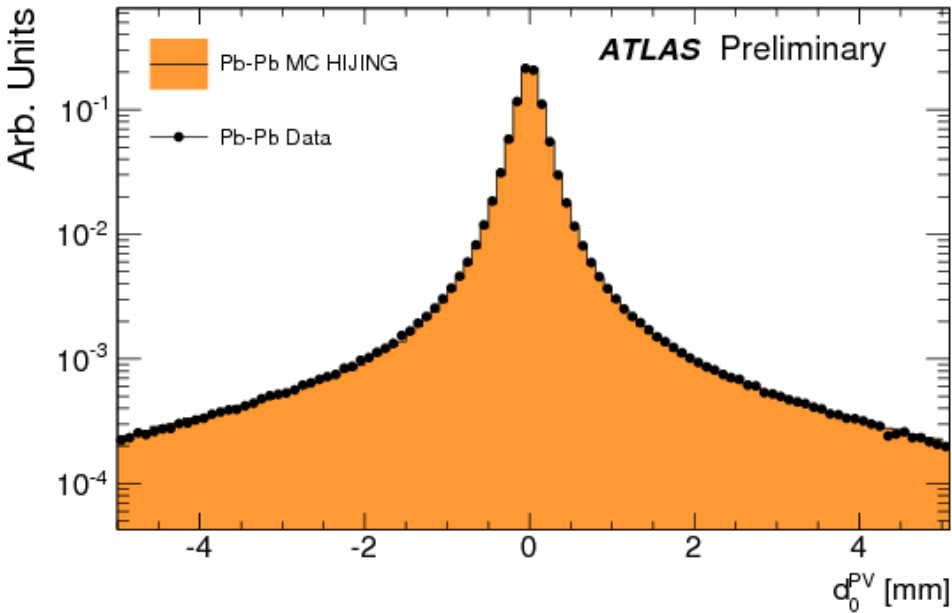
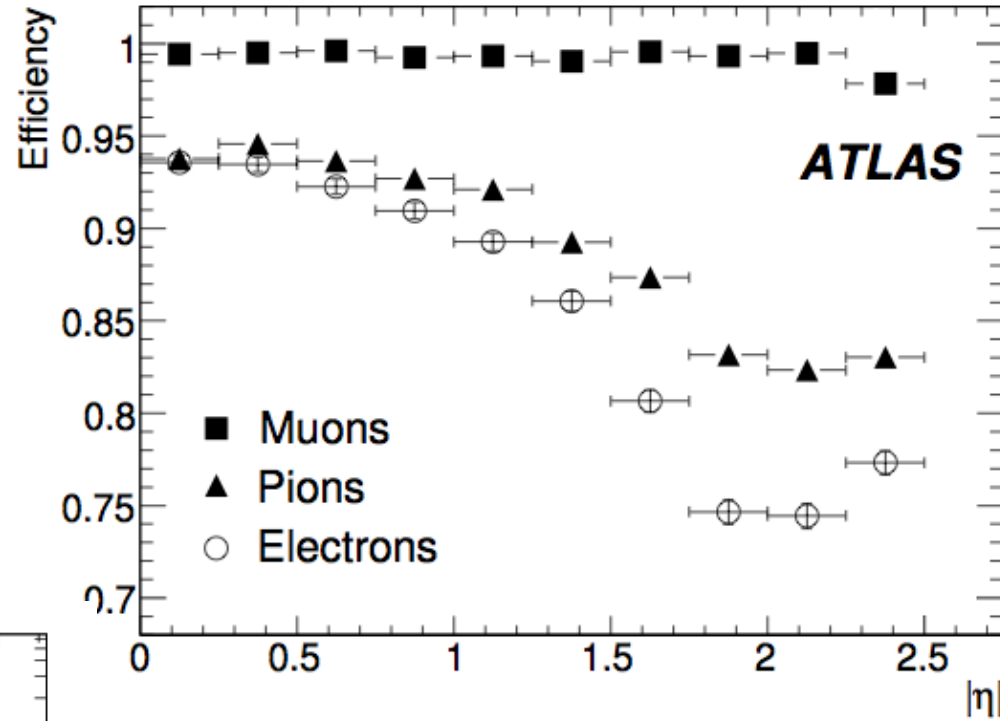
- Pixel detector
  - 50x400 $\mu\text{m}$  pixel size
  - 14x110 $\mu\text{m}$  initial resolution (50 $\mu\text{m}/\sqrt{12}$ )
  - intrinsic resolution down to 3 $\mu\text{m}$  by charge sharing & clustering algorithms
- SCT detector
  - 80 $\mu\text{m}$  strip pitch gives 22 $\mu\text{m}$  resolution
  - stereo angle produces  $\sim 500\mu\text{m}$  resolution in z direction (2<sup>nd</sup> coord.)
- TRT
  - $\sim 150\mu\text{m}$  resolution
- Noisy channels masked, dead channels mapped
  - needs to be known for tracks not to be negatively scored
- Hit efficiency >99% in all 3 systems



# Inner Detector Performance



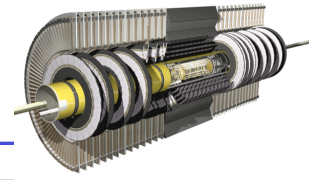
- Track Efficiency: 2 effects
  - detector: loss of track in an interaction  
large effect for electrons, pions
  - algorithms: track finding/fitting  
unsuccessful, small effect
- Measured on Data
  - back-extrapolated muons
  - total number of tracks (mainly pions)



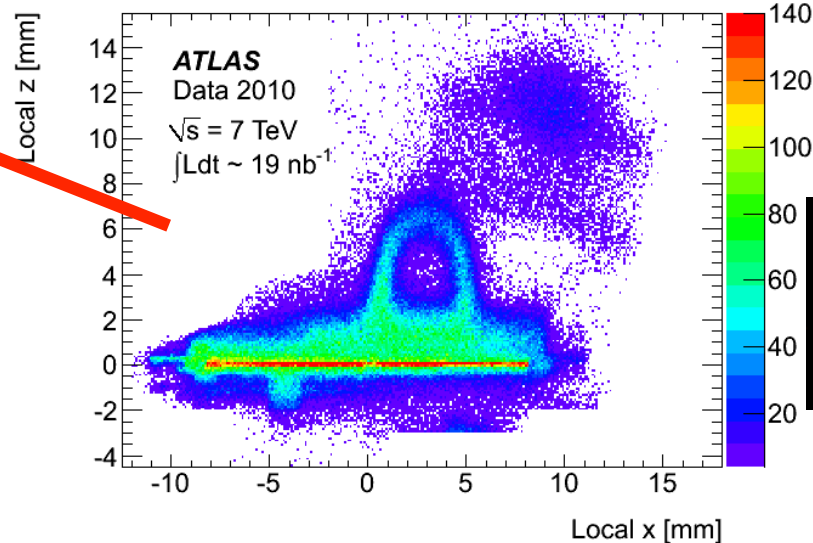
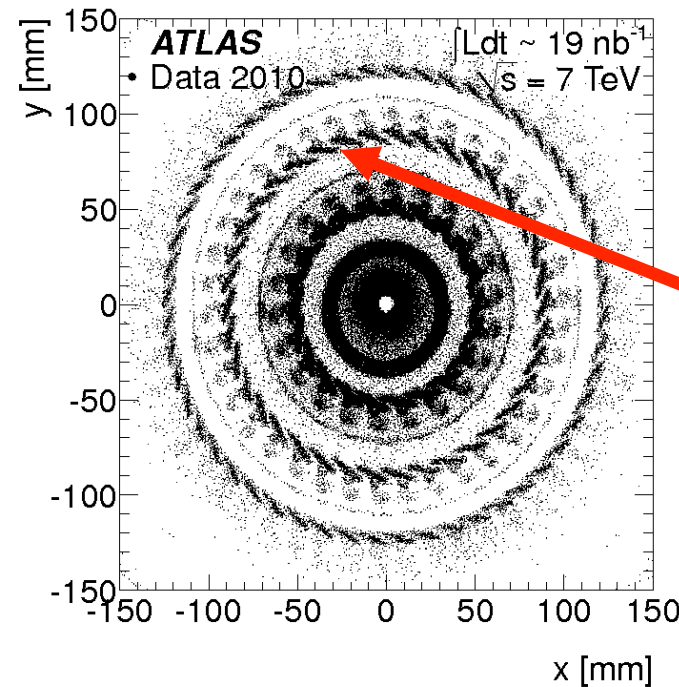
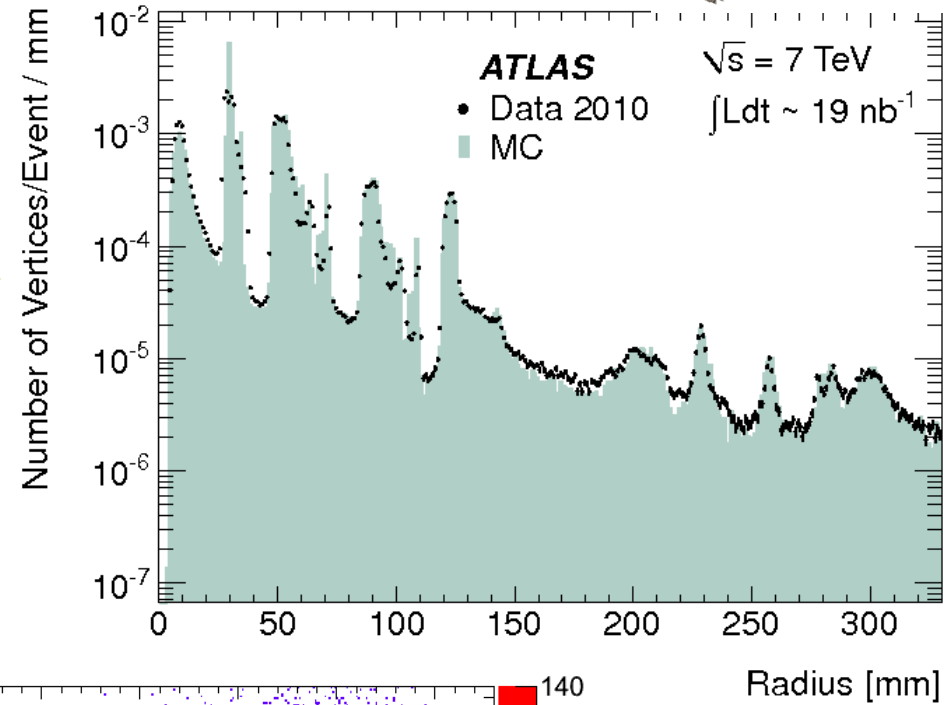
- Impact parameter  $d_0^{PV}$  = track extrapolated to vertex in transverse plane
  - good resolution and good MC model  
crucial for vertexing and b-tagging
  - studied in detail since cosmics,  
here: heavy ion performance



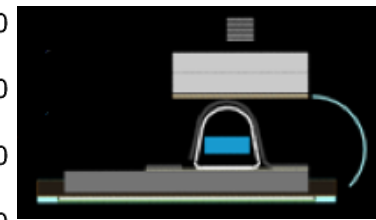
# Material Studies

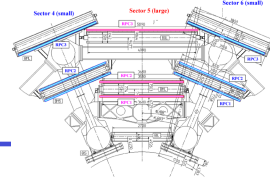


- Interactions in the detector and low mass resonances probe material with high precision
  - $\gamma$  conversion vertex locations
  - hadronic interaction vertices
  - $K^0_s \rightarrow \pi\pi$  and  $J/\psi \rightarrow \mu\mu$  inv. mass
- All were studied, showing that ID material uncertainties < 5%



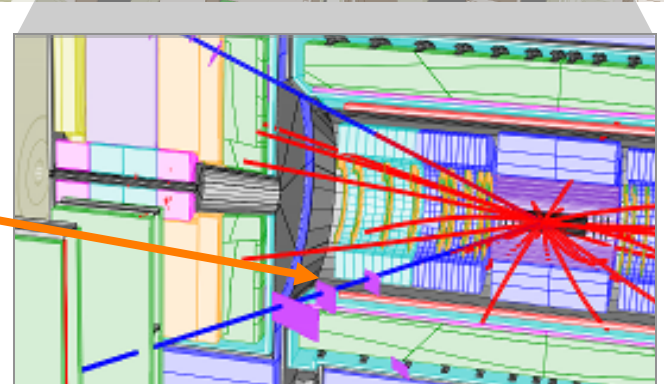
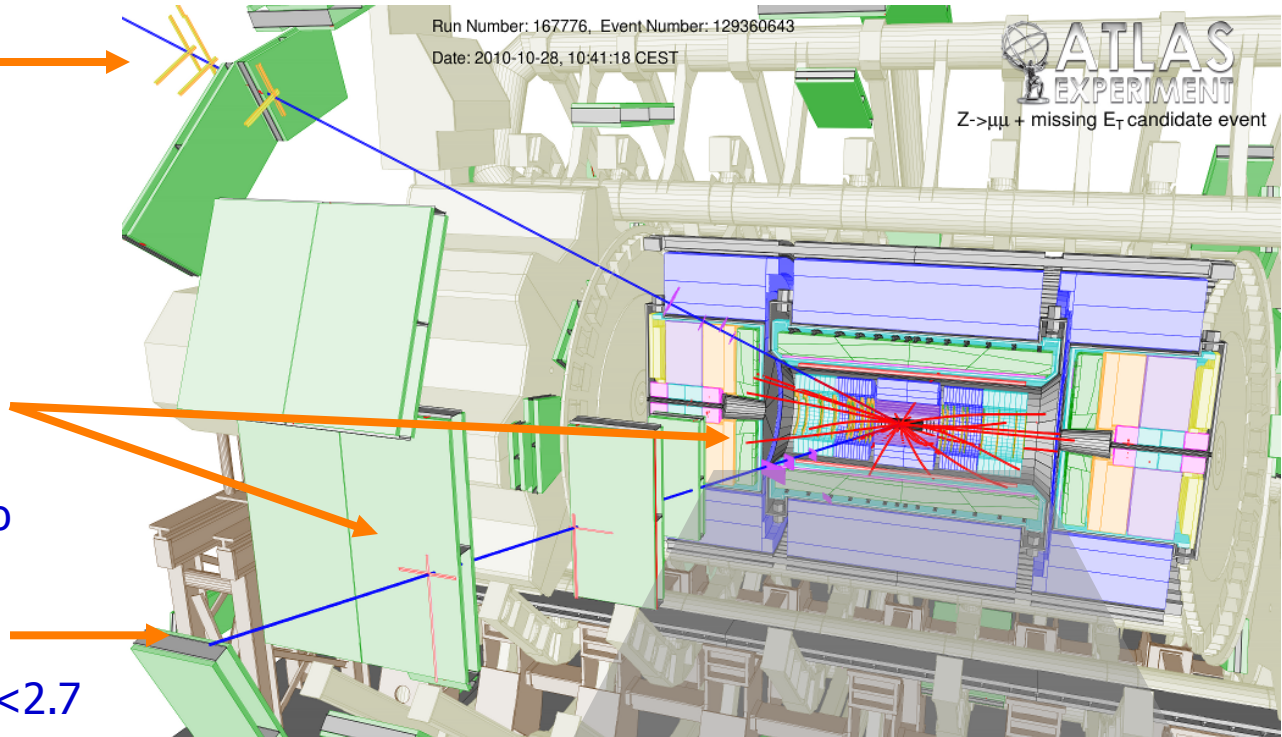
Pixel module in Geant4

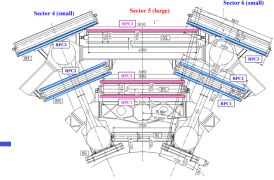




# Muon Reconstruction

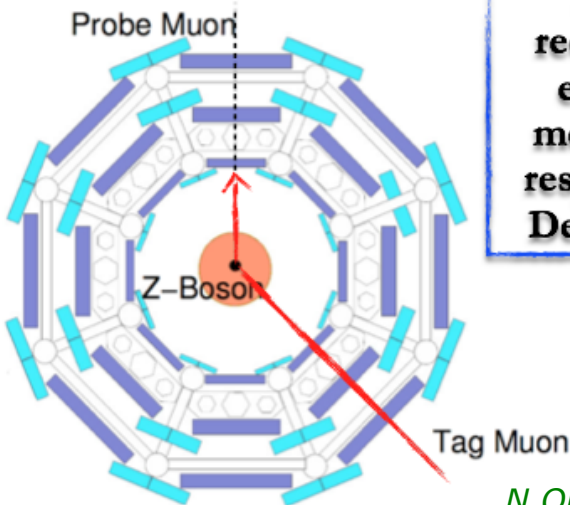
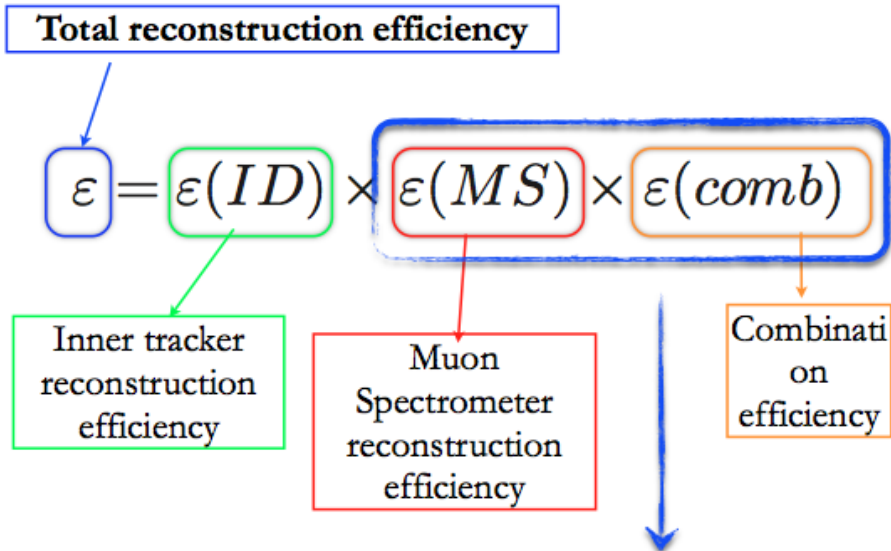
- Segment-tagged
  - available for  $|\eta| < 2.5$
  - most uniform coverage in  $\eta$  and  $p_T$
  - momentum from ID (ID and MS for MuGirl)
- Combined
  - available for  $|\eta| < 2.5$
  - ID and MS contribute to momentum accuracy
- Stand-alone
  - extend coverage to  $|\eta| < 2.7$
  - momentum from MS
  - poor position accuracy at IP
- calorimeter-tagged
  - available for  $|\eta| < 2.5$
  - uniform efficiency near MS acceptance gap at  $\eta \sim 0$
  - optimised for isolated muons



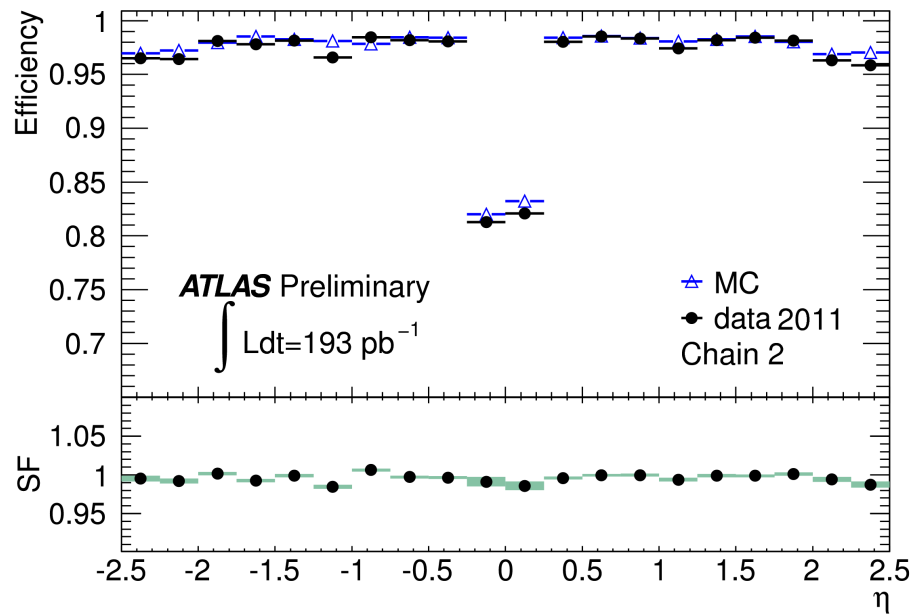


# Muon Efficiencies

- Efficiency from Data: **Tag&Probe**



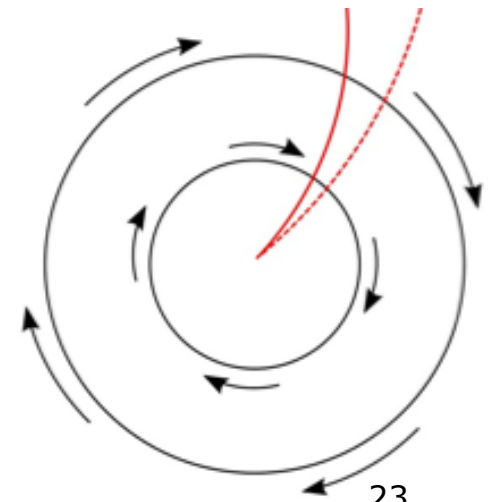
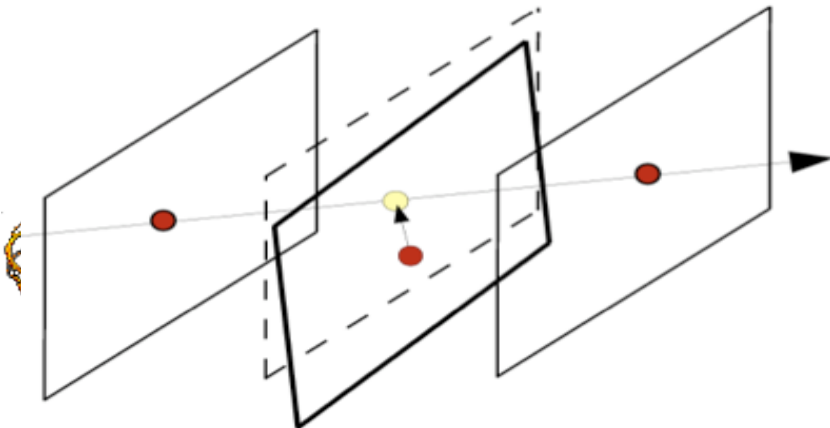
**The muon reconstruction efficiency is measured with respect to Inner Detector tracks**

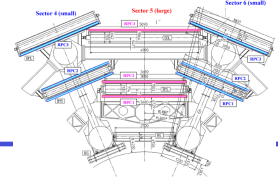


- Results: MC can be used to derive efficiency
  - low pT kinematic range accessible through  $J/\psi \rightarrow \mu\mu$  decay
  - only few cases of scale factor  $SF > \text{error}$

# Alignment of Tracking Detectors

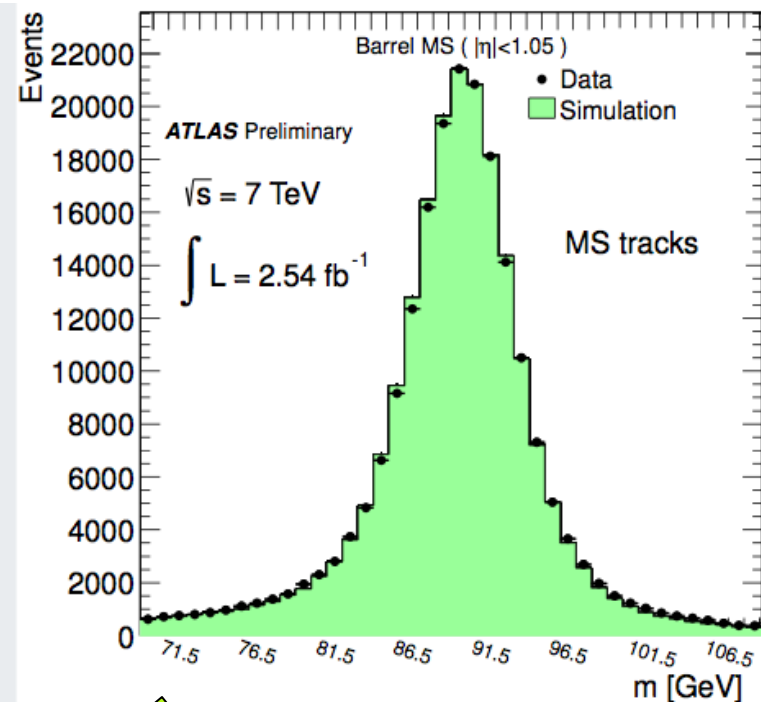
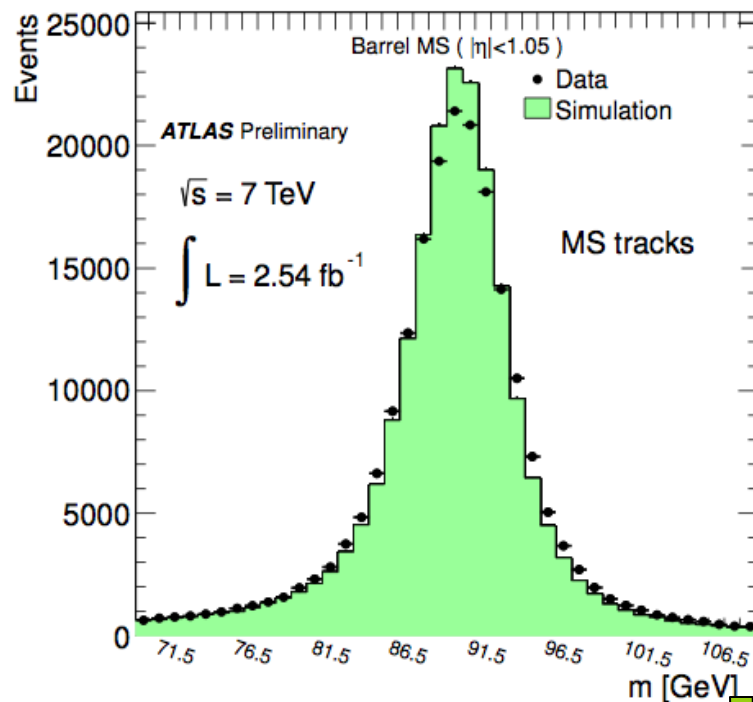
- Detector positioning accuracy
  - $\sim 100\mu\text{m}$  sensors on supports
  - 1-5mm for larger structures
- But: intrinsic resolution 5-100 $\mu\text{m}$
- Positions need to be aligned
  - from data: large track ( $\mu$ ) statistics
  - from detector: optical alignment able to follow “fast” movements
- software alignment is based on minimizing track-hit residuals
  - $X^2 = \sum_{\text{trk}} (r_{TV} - 1r)$ ,  $r = r(\alpha, \lambda_i)$
  - minimization has two big challenges: large # alignment parameters, occurrence of weak modes
- **Large number of parameters**
  - 3 translations + 3 rotations per module
  - ATLAS Pixel+SCT: 5832 modules
- Two algorithmic approaches
  - global chi2: single large matrix including all correlations, fast solving techniques
  - local chi2: solving many small systems, correlations covered by iterations
- **Weak modes**
  - global deformations that do not add to alignment  $\chi^2$  but affect physics quantities
  - curls, twists
  - were studied beforehand but real detector is different





# Muon Momentum Resolutions

- $Z \rightarrow \mu\mu$  decay sensitive to detector effects
  - observed width is superposition of natural decay width and detector resolution
  - $Z$  becomes excellent probe for momentum resolution and scale bias



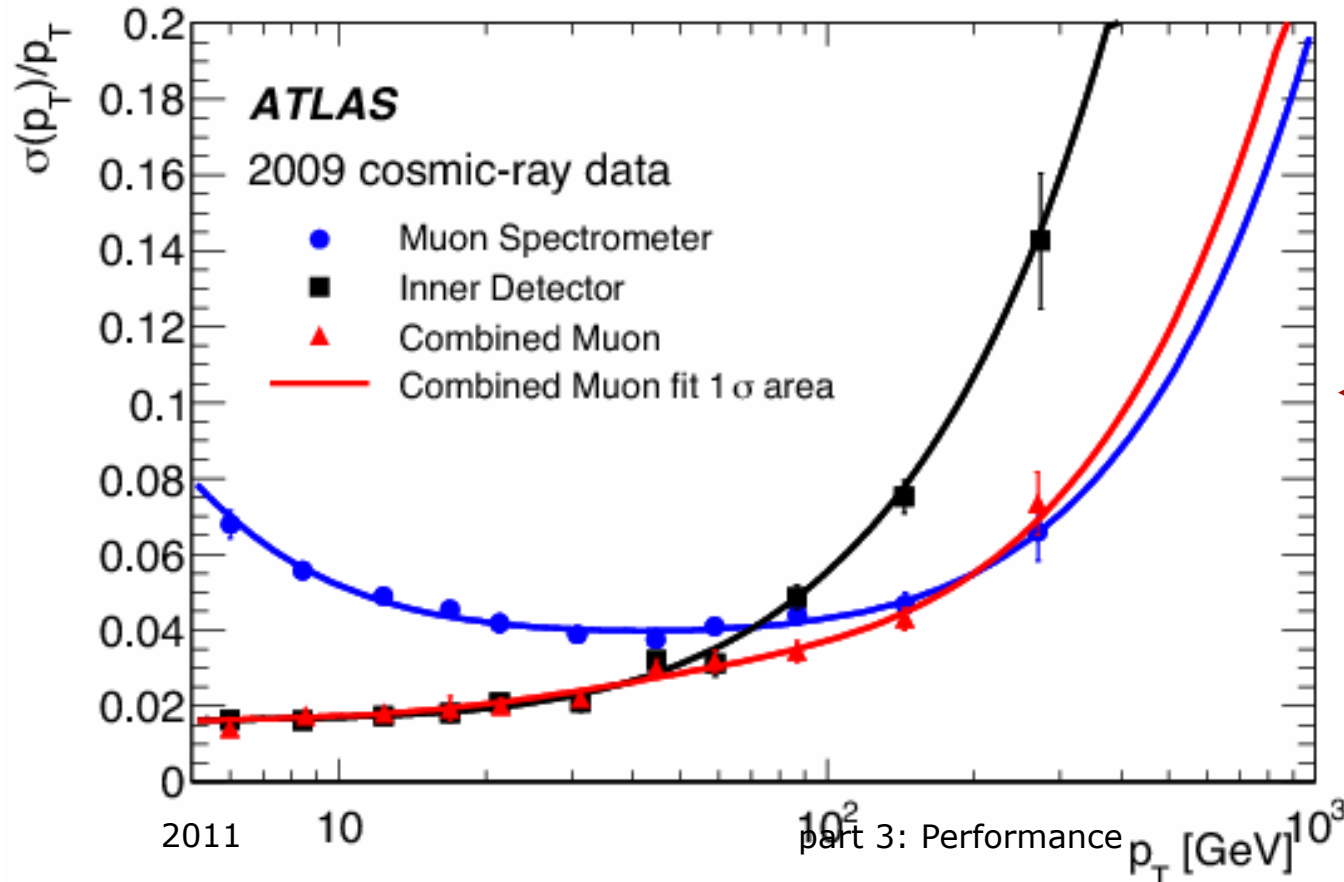
- Momenta on MC smeared before entering physics analyses
  - reduces/quantifies systematic error contribution from cuts, kinematics etc





# Highest $p_T$ Muon Performance

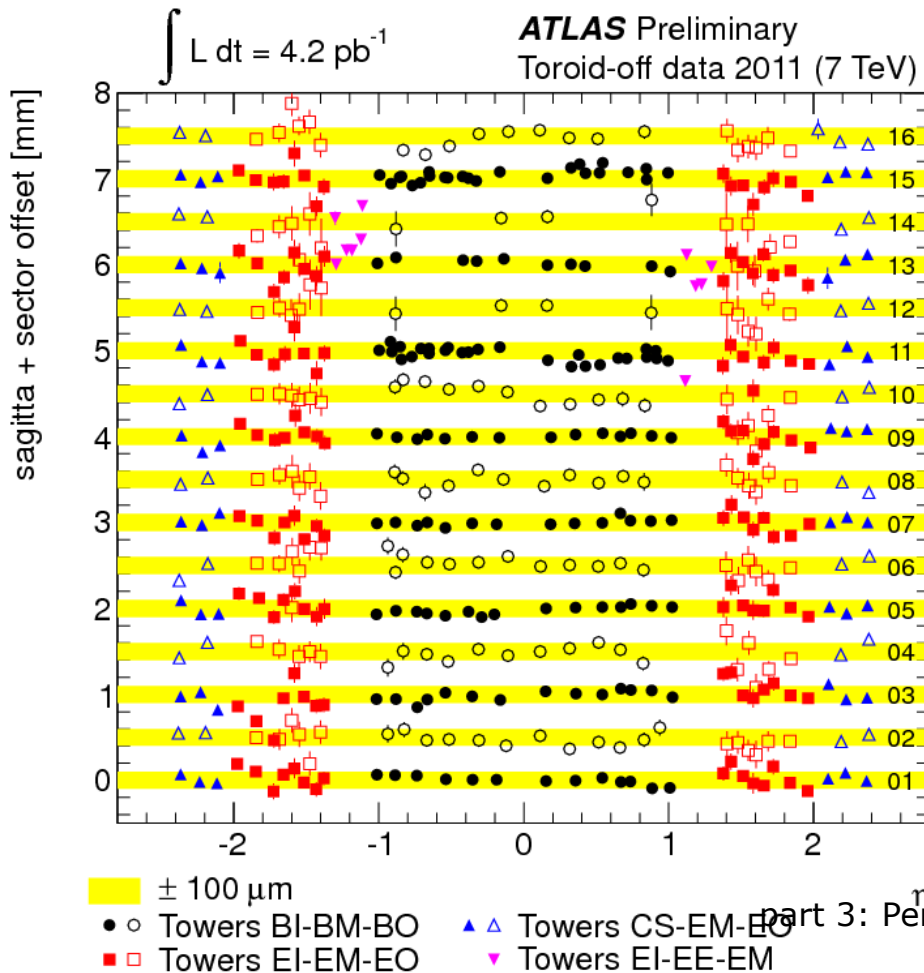
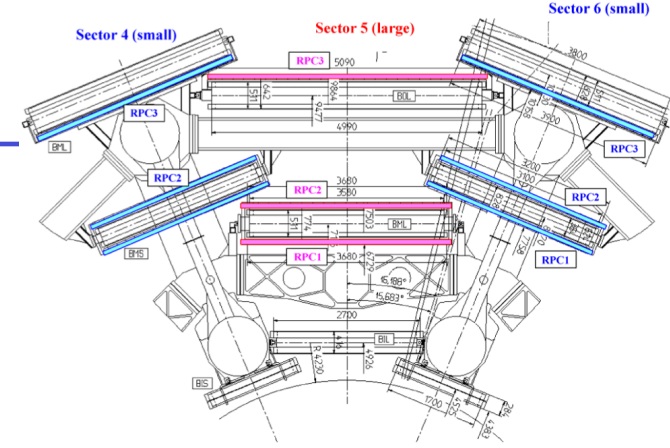
- Muon momenta at  $p \sim 1$  TeV estimated mainly by MS
  - high field integral, lever arm, hit precision in Muon Spectrometer (MS)
  - low momenta determined by ID (material effects in calo+MS strong)
  - TeV scale momentum precision depends on MS alignment
- Precisions alignment of huge MS is a challenge
  - track-based alignment needed to complement and probe quality of optical alignment



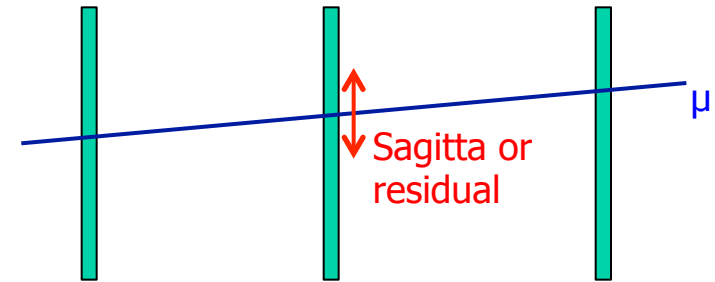
← Design performance is 10% at 1 TeV

# High $p_T$ Performance

- Methods to achieve precision alignment
  - special runs: solenoid off, toroid on
  - overlaps between station sectors in  $\Phi$
  - cosmic rays (mainly 2009 preparation, see previous slide)



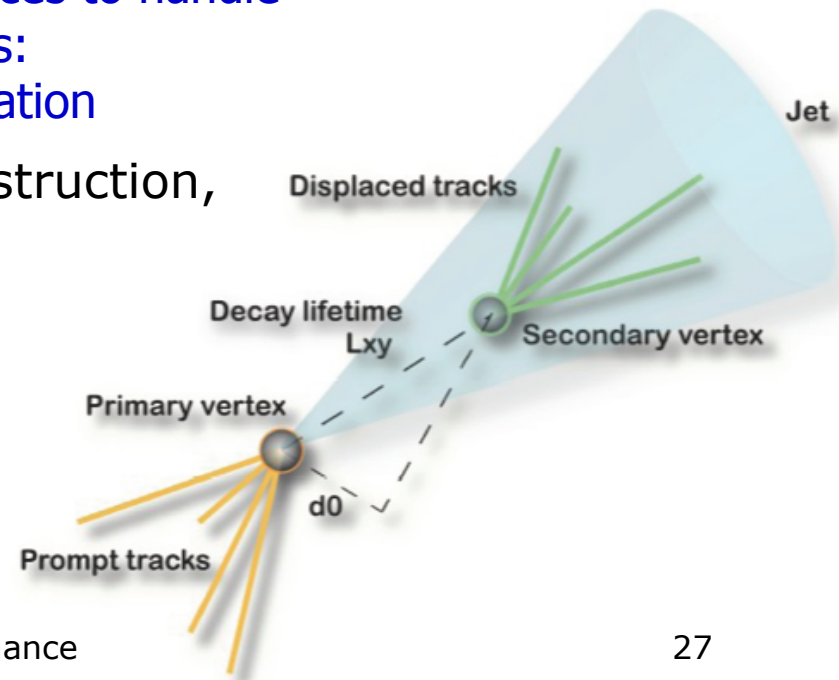
- Solenoid-off runs:
  - track muons using momentum from ID (mat. effects correctly parameterized)
  - study sagitta in 3-point system: should be 0 for straight tracks



- difference from 0 allows alignment
- optical alignment follows movements when toroids switched back on (!)
- typical precision now  $< \sim 100 \mu\text{m}$

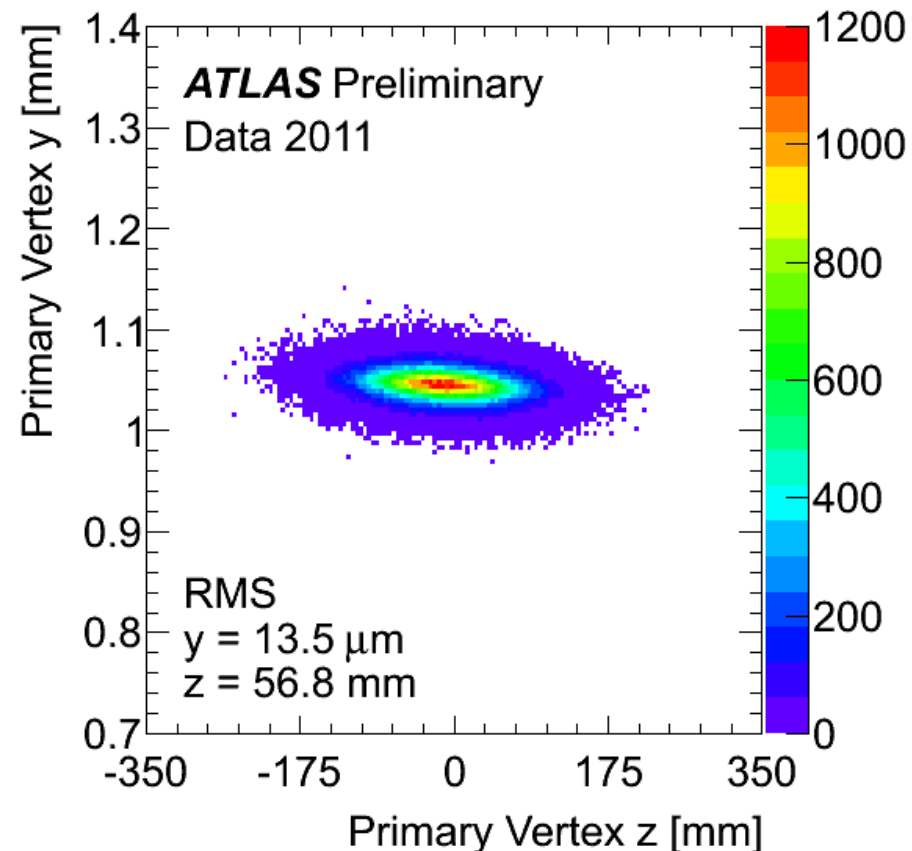
# Vertex Reconstruction

- Task is to estimate vertex position  $\nu$ 
  - from N track parameter vectors (& cov)
  - and update track momentum vectors and full  $3+3N$  covariance matrix
- Algorithms are similar to track fit
  - measurement model  $\lambda_i = h(\nu, p_i)$  describes how track parameters depend on vertex and momenta at vertex
  - model is inherently not linear, large matrices to handle
  - use again linear estimators with iterations: progressive (Kalman) and global minimization
- Applications are primary vertex reconstruction, heavy flavour decay vertices, b-tagging,  $\gamma$  conversions, kinematic fits ...



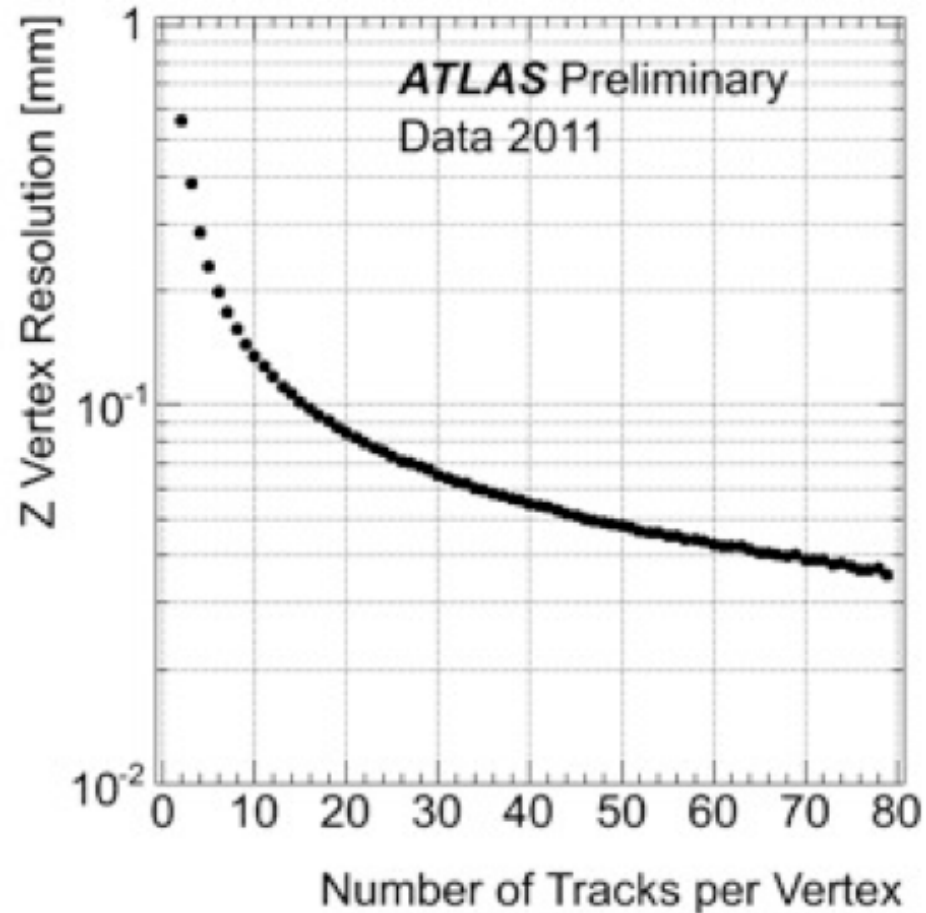
# Vertex Finding

- Primary vertexing at LHC combined finding and fitting **adaptive multi-vertex fitters**
  - iterative, reweighted Kalman filter
    - Kalman filter = adds tracks progressively to vertex candidate
  - robust fitter: outlying tracks are down-weighted automatically
  - new vertex candidate formed with outlying tracks (minimally 2 tracks form vertex)
  - list of vertex candidates is input to next iteration, vertices compete against each other for tracks
- **Beam spot**
  - cloud of primary vertices averaged over short period in time
  - routinely determined in data-taking
  - beam spot then used as constraint in primary vertex finding
- Measures **in-time pile-up**
  - $\mu$  = number of pile-up vertices



# Vertex Reconstruction Perf

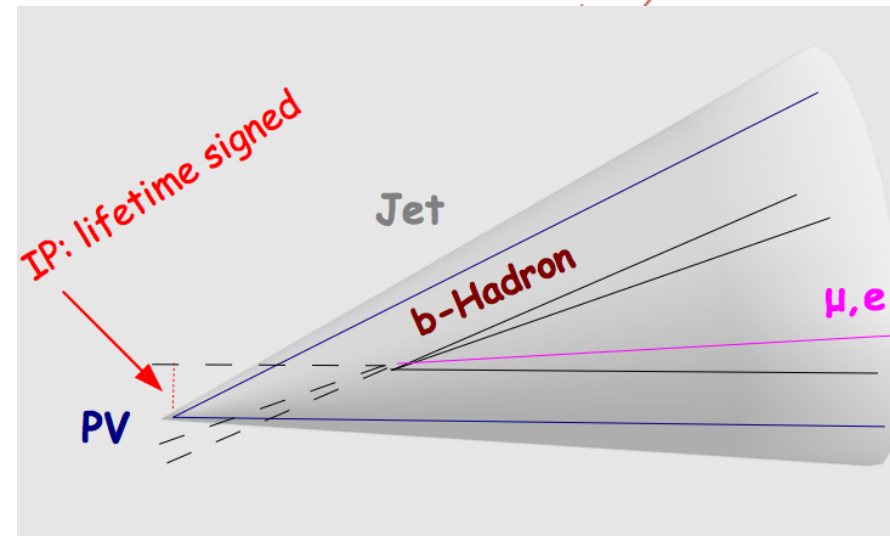
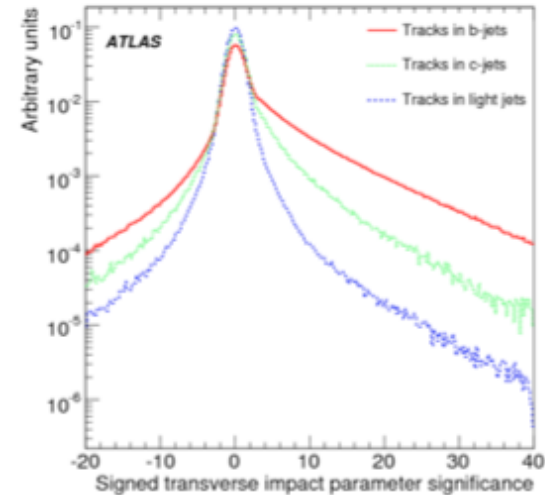
- Primary vertex in another physics object for which we need to know precisely its performance
- Vertex resolution on data measured with split-vertex method
  - randomly split PV into two
  - study difference between positions from two vertex fits
  - expect 0 with variance from resolution
- Resolution depends on number of tracks at vertex





# b-Jet Tagging

- ▶ Spatial tagging (or life-time tagging):
  - B hadrons have a significant flight path length:
    - $E(B) \sim 50 \text{ GeV} \Rightarrow L \sim 5 \text{ mm}$
  - Secondary vertex in jets.
  - Tracks with high positive impact parameter.
  
- ▶ Soft lepton tagging: Useful to commission other taggers
  - Low  $p_T$  electron/muon from B/D decay.
  - Efficiency limited by (B/D  $\rightarrow$  l) branching ratio.



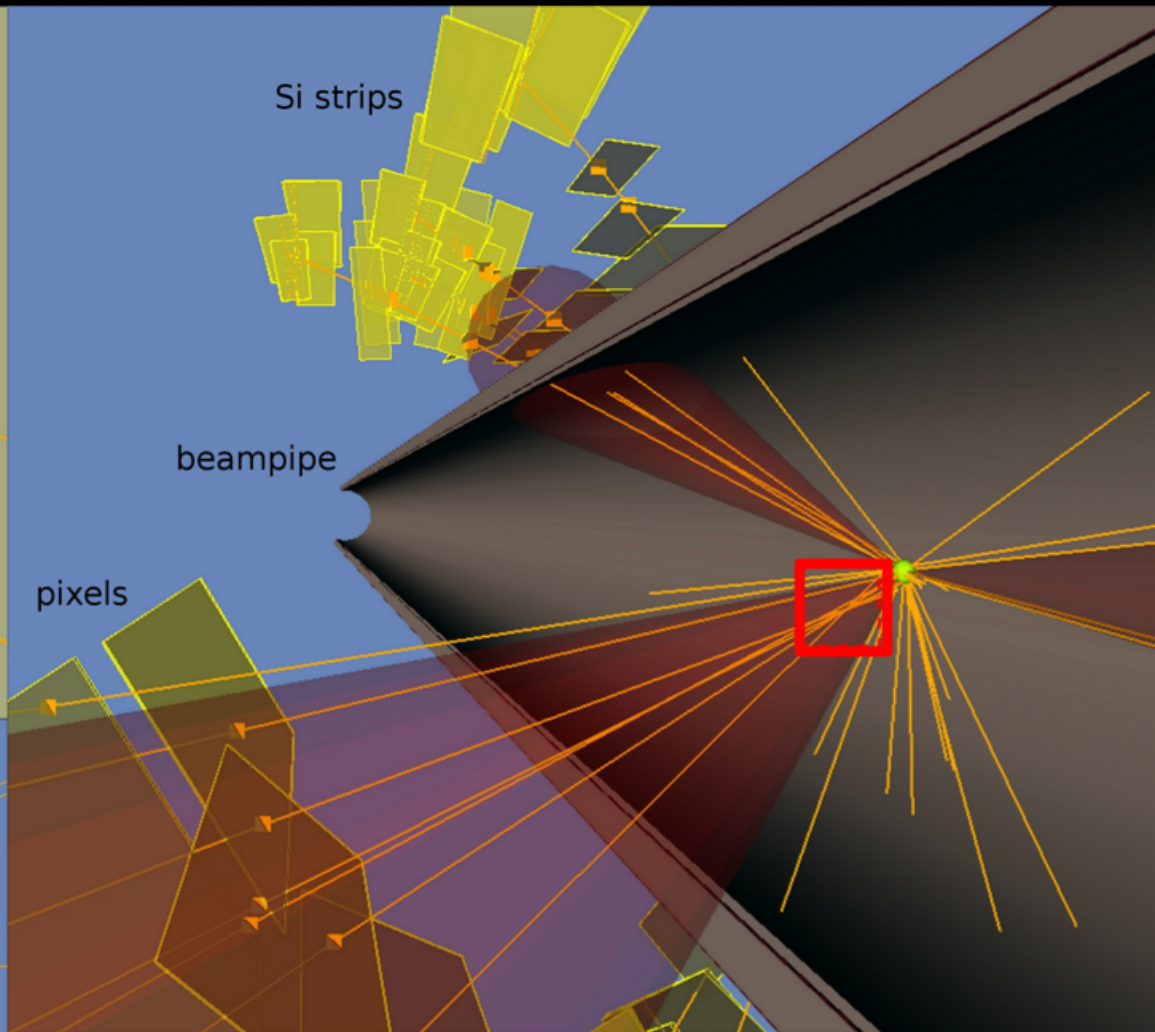
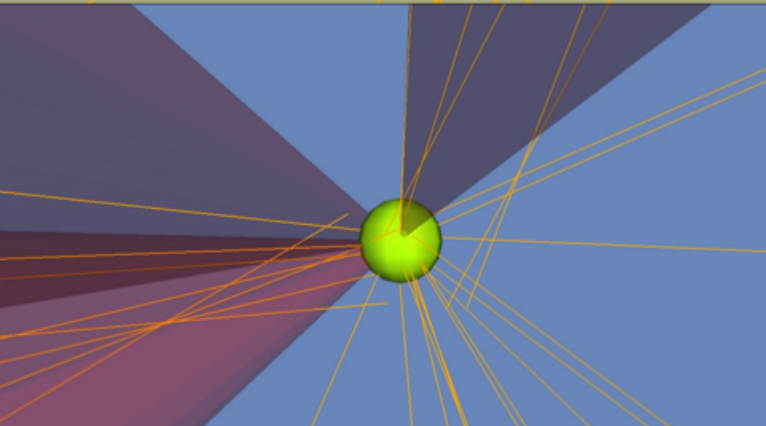
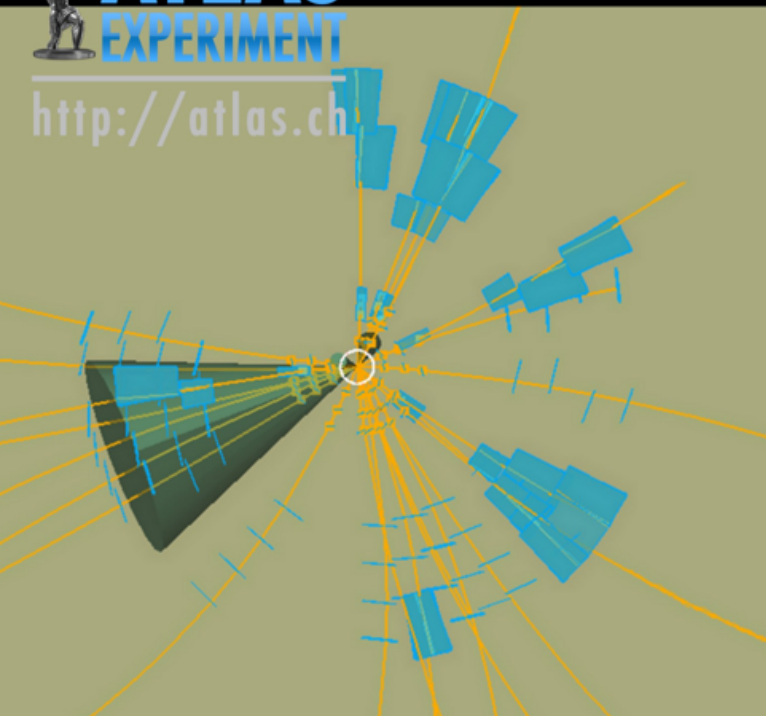
# b-tagged jet



Run 152166  
Event 817271

<http://atlas.ch>

b-tagged jet in 7 TeV collisions



jet  
 $p_T = 19$  GeV (measured at electromagnetic scale)

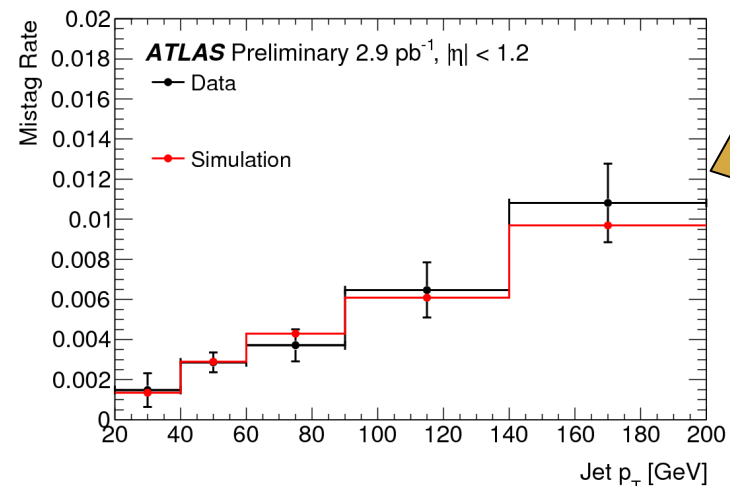
# X B-Tagging Inventory

- ▶ **Simple taggers:** *robust, used on initial data (2010)*
  - **TrackCounting:** Counts tracks with high IP
  - **JetProb:** Track compatibility with the primary vertex
  - **SV0:** flight length significance of the SV
- ▶ **Advanced taggers :** *After commissioning, used for 2011 physics results*
  - **IP $n$ D** ( $n=1,2,3$ ) : IP based likelihood tagger
  - **SV $n$**  ( $n=1,2$ ): SV based likelihood tagger
  - **JetFitter $X$**  ( $X=$ Tag,TagNN,COMB,COMBNN)
- ▶ **Soft lepton taggers :** *Limited efficiency, also tool for calibration*
  - SoftMuonTag
  - SoftElectronTag



2011

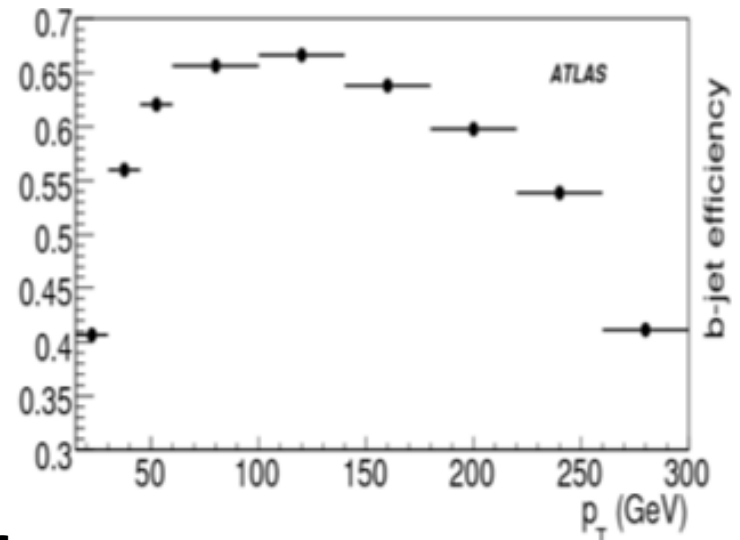
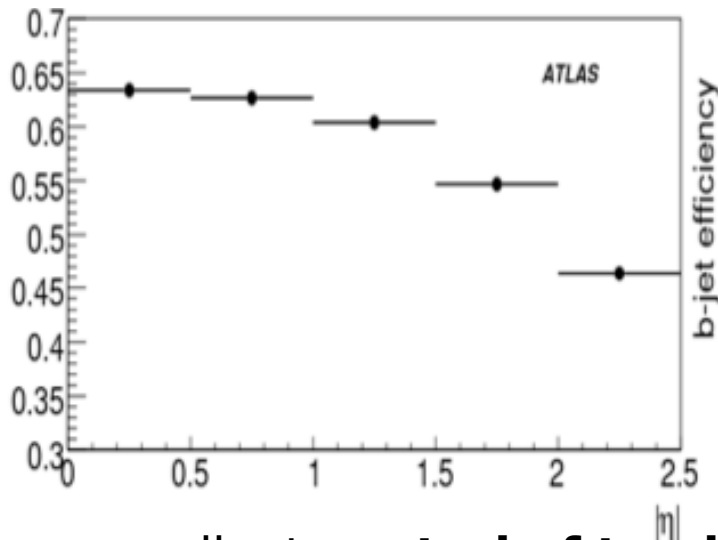
part 3: Pt



32

# B-Tagging Challenges

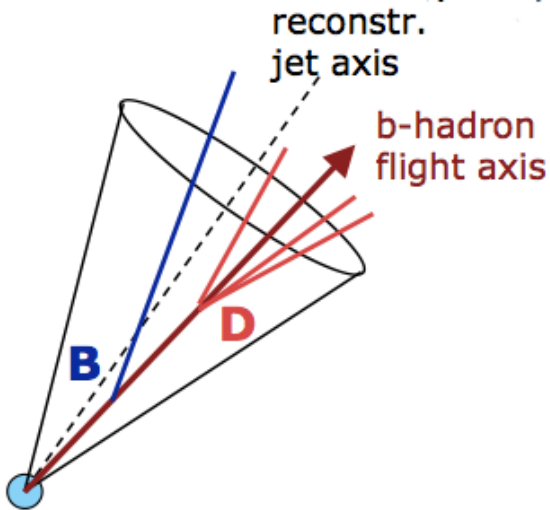
- strong dependence on **kinematics**
  - low  $p_T$  and high  $|\eta|$ : multiple scattering and material interactions
  - high  $p_T$ : two effects to cope with:
    1. collimated tracks  $\rightarrow$  limits of pattern recognition
    2. 'late' B decays in detector ( $p_T \sim 200$  GeV: 8% decay after b-layer)
  - shown for  $t\bar{t}$  events: efficiency of IP3D+SV1 tagger at cut  $w > 4$



- excellent **control of tracking performance**
  - good local alignment, material description
  - study of impact parameter and vertex resolutions
- Use of sophisticated taggers in 2011 results is consequence of excellent tracking performance and collaboration tracking-btagging



# JetFitter

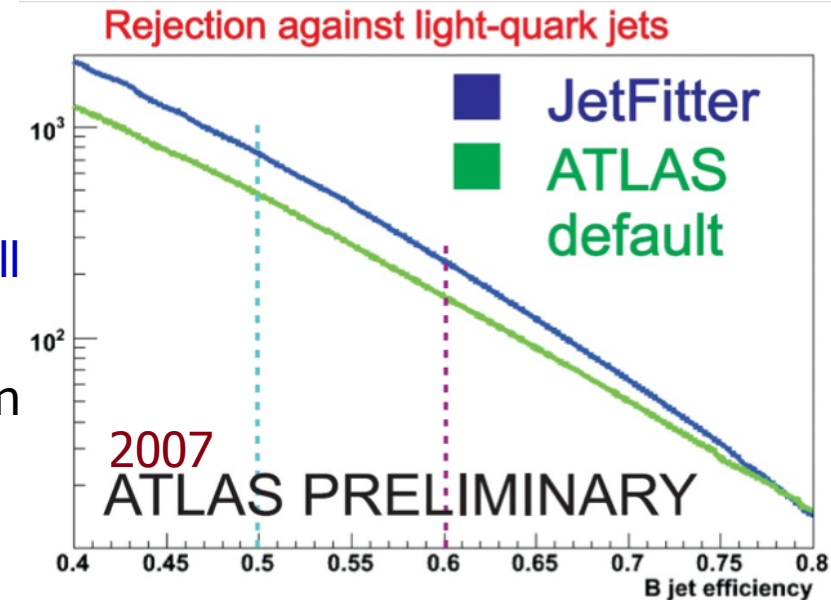


- B decay is actually a decay chain:
  - $B \rightarrow D \rightarrow K/\pi$  with significant D decay length
- SV taggers improve b-tagging but do not use optimal/accurate information
  - $B \rightarrow D$  cascade approximated by single vertex (# tracks and resolution not enough to fit 2<sup>nd</sup> + 3<sup>rd</sup>)
  - contaminated with light flavour jet ( $K_0$ ,  $\Lambda$  decay)
  - statistical issues with 1-vertex assumption ( $\chi^2$ , cov)

- JetFitter algorithm solves issues with a multivertex fit in 1 dimension along the jet axis: PV – B-vertex – D-vertex
  - robust against small number of tracks (1)
  - displacements from common jet axis small
  - Kalman filter based
- B hadron discriminators extracted from B-D system ( $m$ ,  $E/E_{jet}$ ,  $\sigma(d)/d$ )



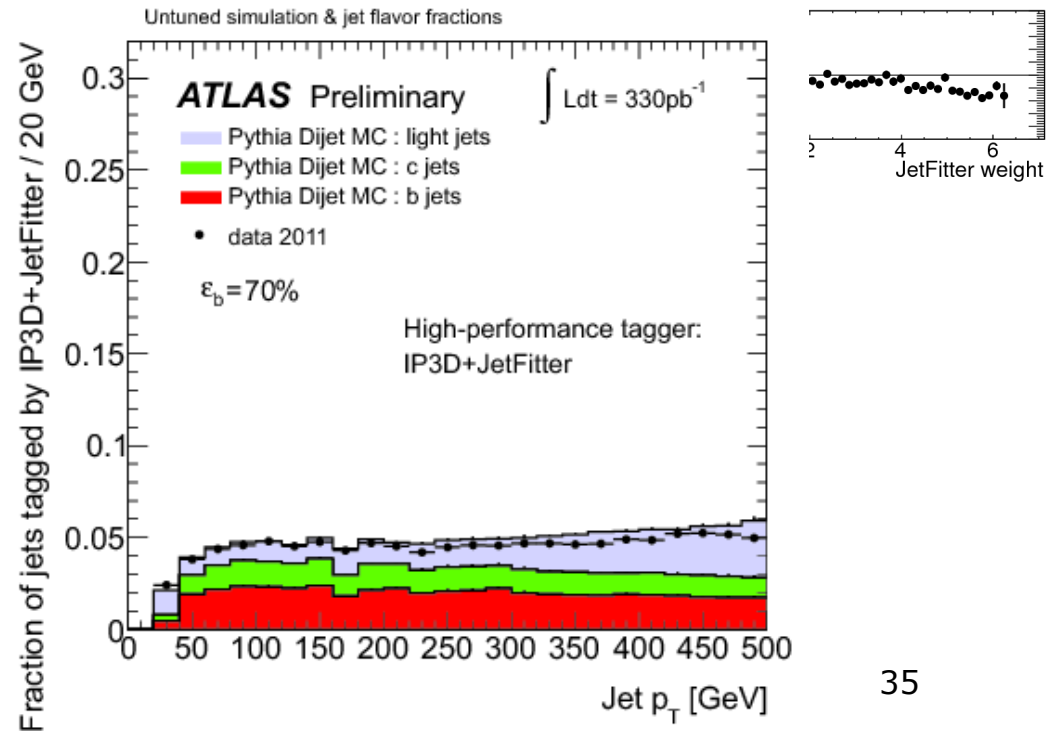
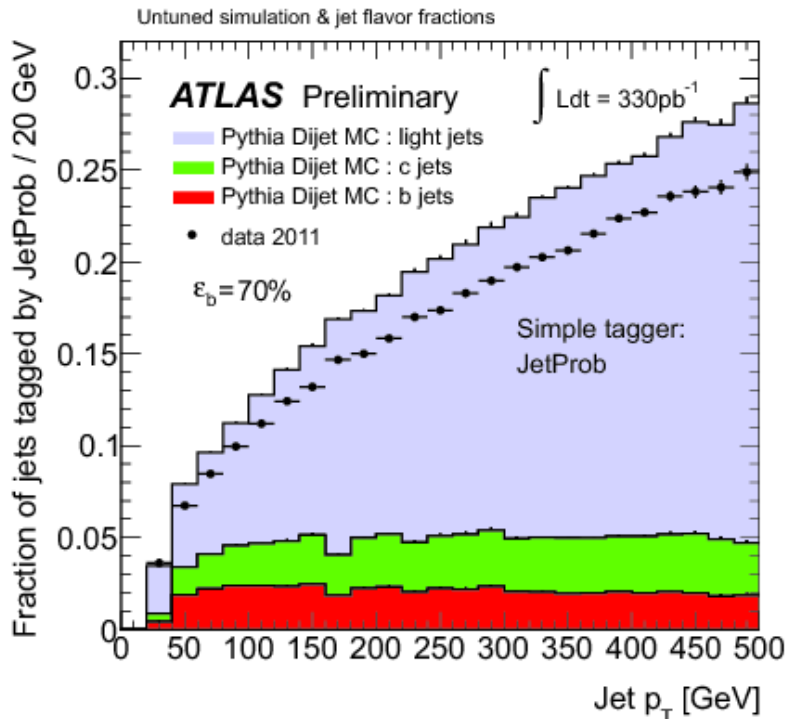
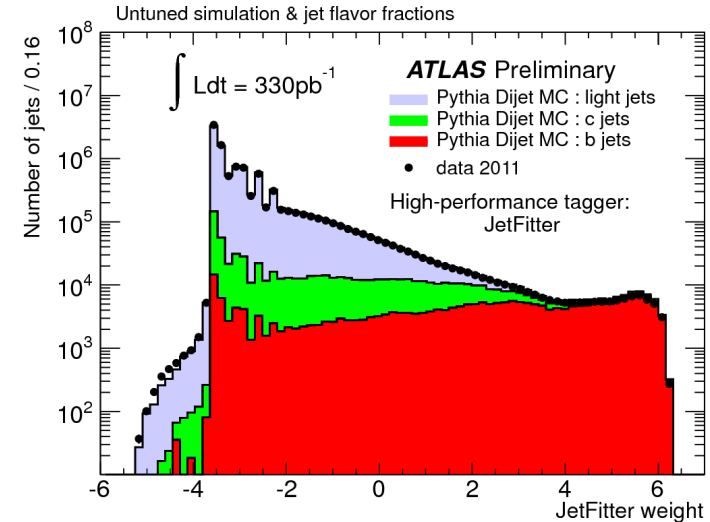
indeed rejection of light jets improved





# b-Tagging Performance

- Comprehensive studies of all b-taggers
  - input variables and output weights
  - relative comparisons in different kinematic regions
  - efficiencies
- Powerful combination IP3D+JetFitter



# Electrons, Jets, Missing $E_T$ , Tau



# Electron Identification

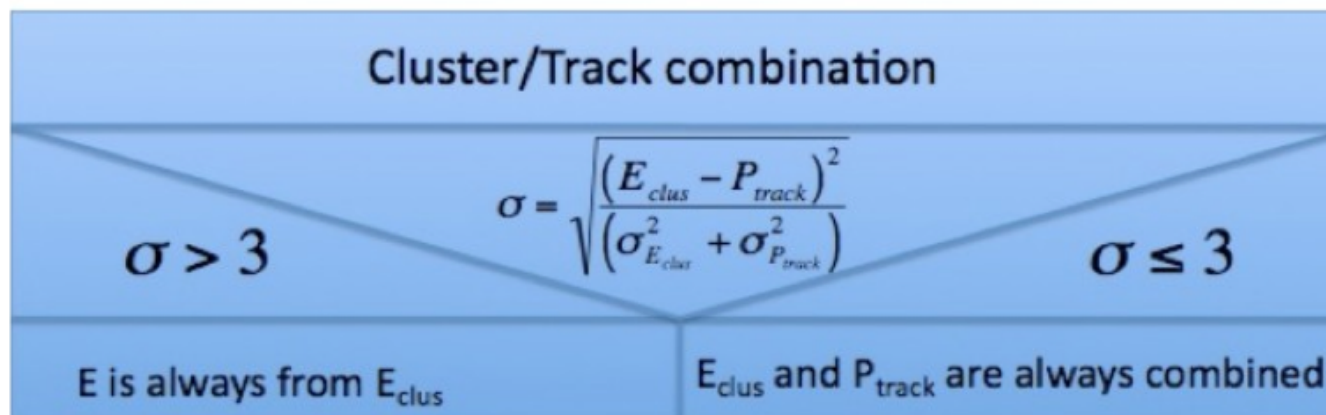
- standard identification: LAr el-mag. **calorimeter seeded**
  - seeded by clusters reconstructed in LAr by a sliding window algorithm
  - attempt to match a track to the cluster
  - attempt to match a conversion vertex to the cluster
- Definition of objects
  - electrons: cluster + track
  - photons: two categories:
    1. unconverted photon
      - = cluster + no track, no conversion vertex
    2. converted photon = cluster + conversion vertex
- additional identification: **track seeded**
  - tight pre-selection cuts to minimize false identification
  - keep standard track+cluster if track is the same
  - improves efficiency at low energies,  $ET < 5\text{GeV}$
- **forward electrons**
  - uses topological clusters, no InDet information  $|\eta| > 2.5$
  - dedicated identification algorithm



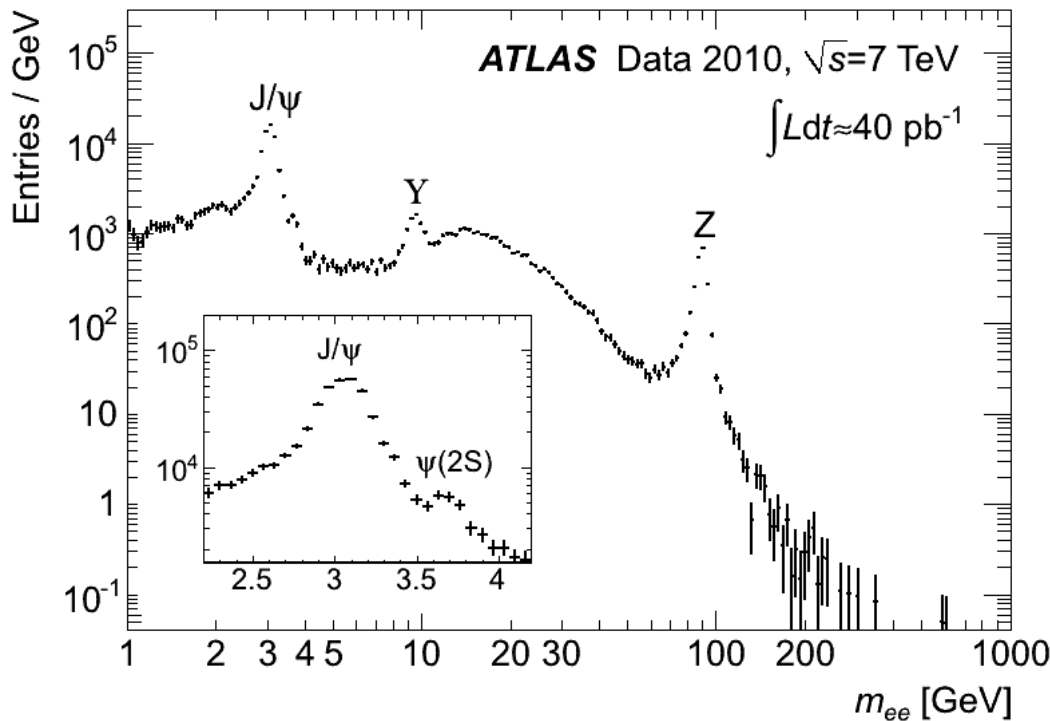
# Electron Energy Reconstruction

egamma objects are massless, with four-momentum defined:

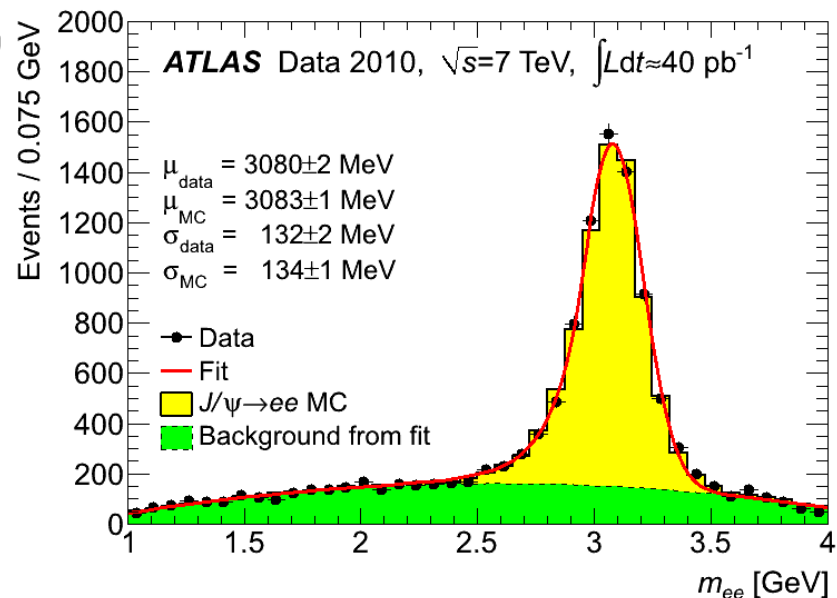
- ★ For electrons: if  $\sigma < 3$  and track is not low  $p_T$  TRT-only, the energy is from combining the cluster energy and the track momentum; else it comes from the cluster.  $\phi$  is from the track, and  $\eta$  comes from the track, unless the track is TRT-only, in which case, the  $\eta$  is from the cluster pointing
- ★ For unconverted photons, energy is from the cluster.  $\eta$  comes from cluster pointing, and  $\phi$  is from the cluster position. From 15.8,  $\phi$  is corrected for the primary vertex.
- ★ For converted photons, energy is from the cluster,  $\phi$  is from the track, and  $\eta$  comes from propagating from cluster to conversion vertex, unless the tracks are TRT-only, then cluster pointing



# Electron Performance



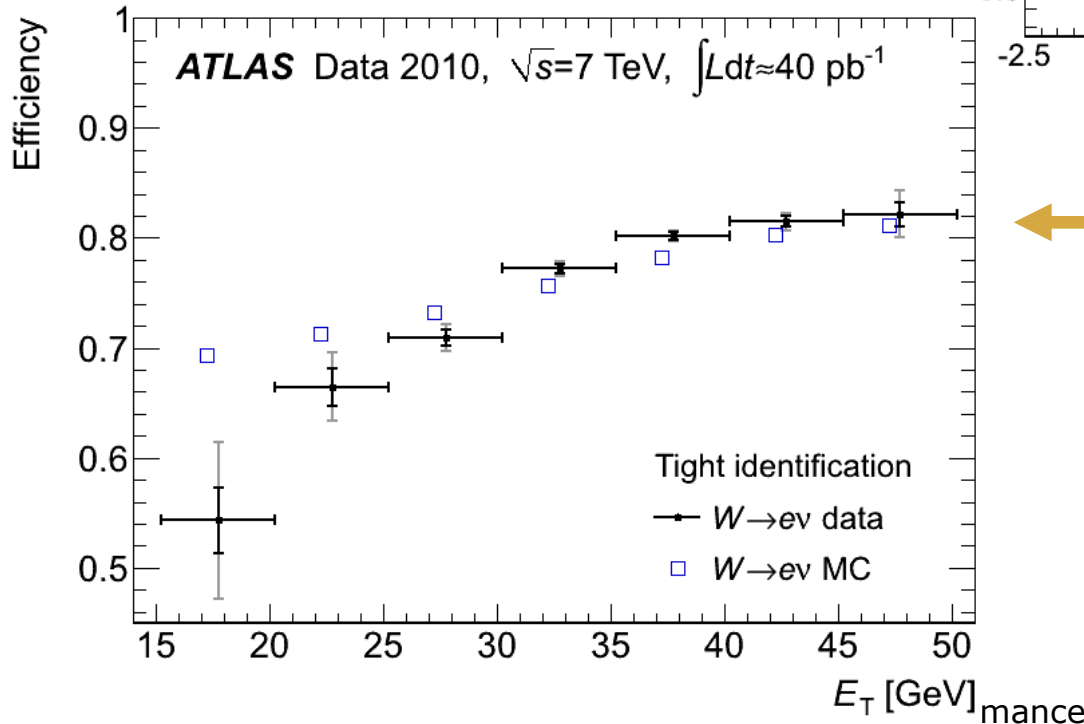
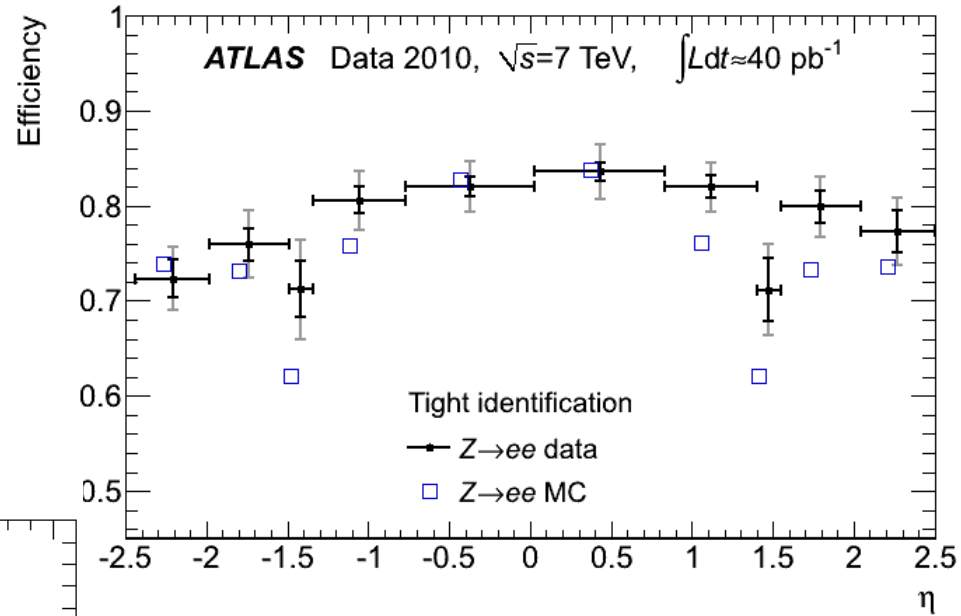
- same di-lepton spectrum as for muons
- momentum resolution at low  $p_T$ : wider peaks or tighter selection
- well described by MC





# Electron Performance

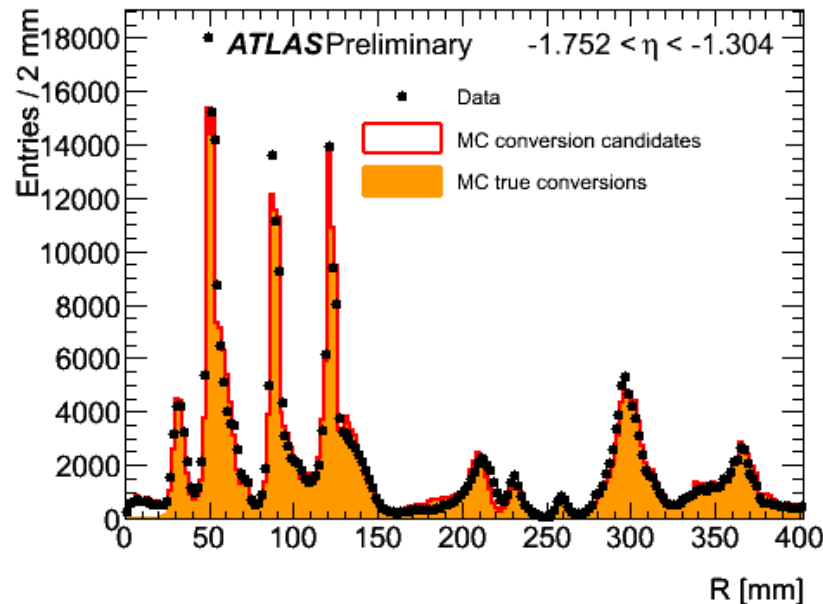
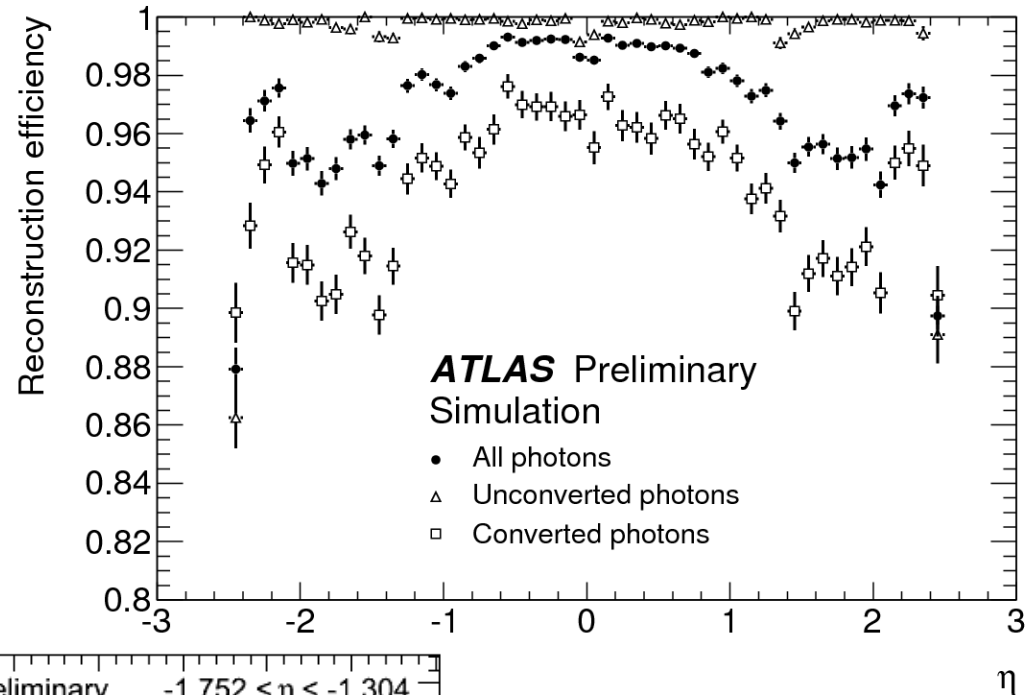
- efficiency from  $Z \rightarrow ee$ 
  - tag-and-probe like with  $\mu$
  - tight identification = few false ID
  - note the low efficiency  $\sim 80\%$  compared to muons



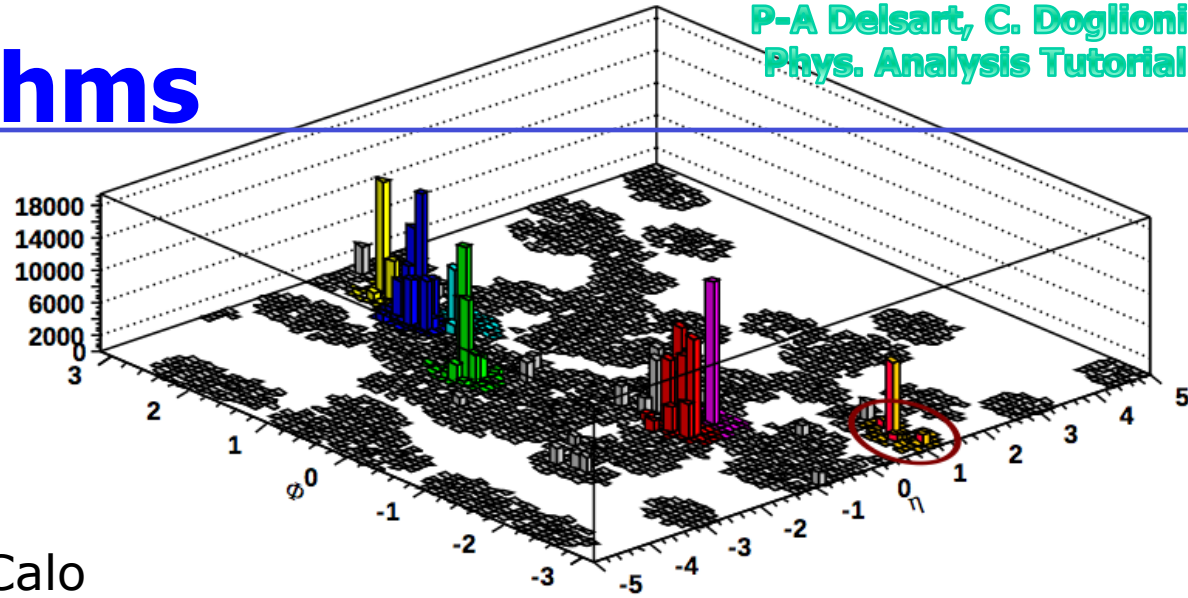
- efficiency from  $W \rightarrow e\nu$ 
  - reaches lower in  $E_T$
  - some discrepancies (note: 1 year old reco software)

# Photon Performance

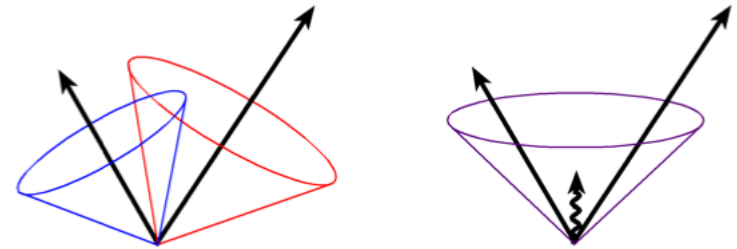
- more difficult to estimate
- conversions are main issue
  - main contributor to inefficiency, well known from simulation
  - studies therefore focus on understanding the material
  - conversions complement other material studies (slide 20)



# Jets Algorithms



- Task: estimate direction and energy of prompt hadrons from energy map in Calo
- Geometrical cone algorithms simple but not infrared safe →
- Kt/Cambridge clustering algorithms
  - define distance and limit



$$d_{ij} = \min(p_{Ti}^2, p_{Tj}^2) \frac{\Delta R_{ij}^2}{R} \quad \text{and} \quad d_{iBeam} = p_{Ti}^2$$

- keep merging two smallest distance objects  $i, j$  into new proto-jet until  $d_{ij} > d_{iBeam}$



# Jet Clustering Algorithms

Kt algorithm is collinear and Infrared safe. But has inconvenient :

- Irregular, complex shape
- “vacuum cleaner” effect

Some variants have been studied : replace distance definition by

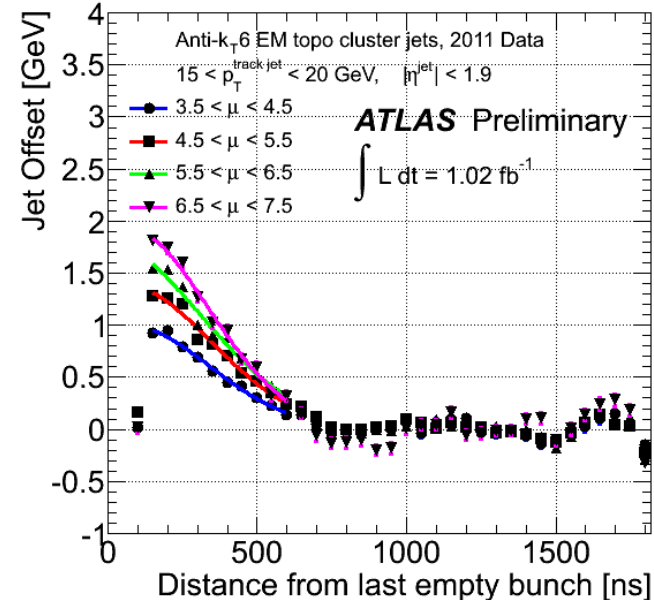
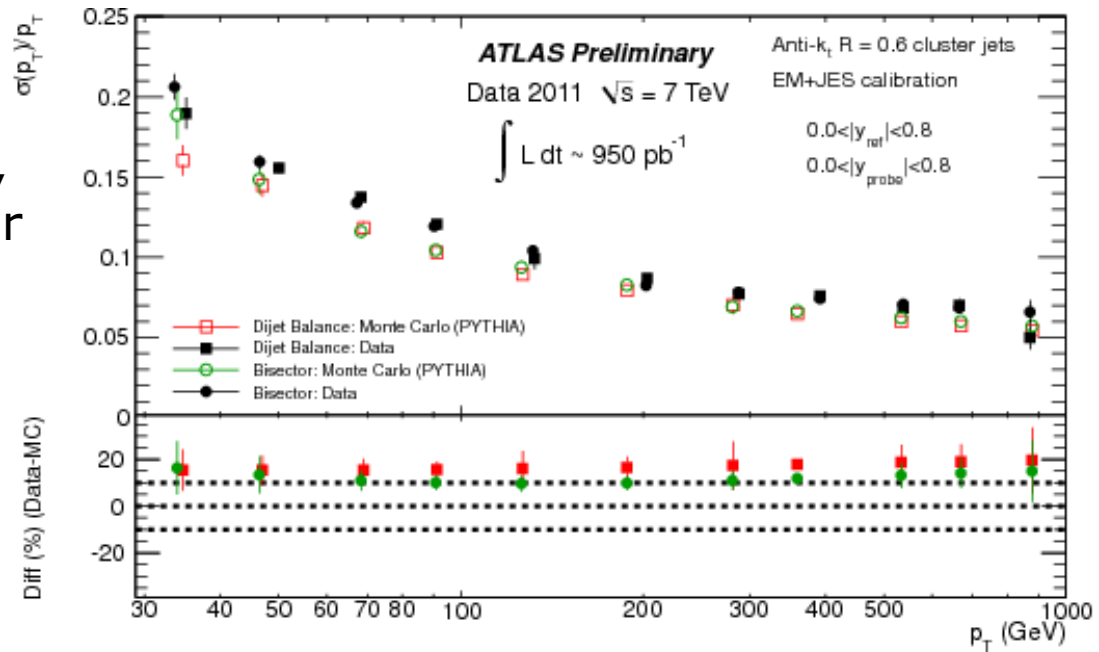
$$d_{ij} = \min(p_{Ti}^{2p}, p_{Tj}^{2p}) \frac{\Delta R_{ij}^2}{R^2} \quad \text{and} \quad d_{iBeam} = p_{Ti}^{2p}$$

- 1 p=1 : Kt algorithm. Priority to low Pt constituents
- 2 p=0 : Cambridge variant. Purely geometrical
- 3 p=-1 : **Anti-kt** variant. Priority to high Pt constituents
- 4

**Anti-kt** very recent (Salam, Cacciari, Soyez [arXiv:0802.1189](https://arxiv.org/abs/0802.1189), independently of Atlas development [Atlas CVS ;)), has several advantages.

# Jet Reconstruction Performance

- Mainly energy resolution and reliability of jet energy, the **jet energy scale** factor (JES or JSF)
- various in-situ methods
  - in situ = measure on data
- look at known balanced events, like di-jet
- Another effect:
  - **out-of-time pile-up**
  - Calo sensitive to energy from preceding collision
  - Energy may be overestimated
  - study by comparing to track-jets (hits on track have tighter timing)



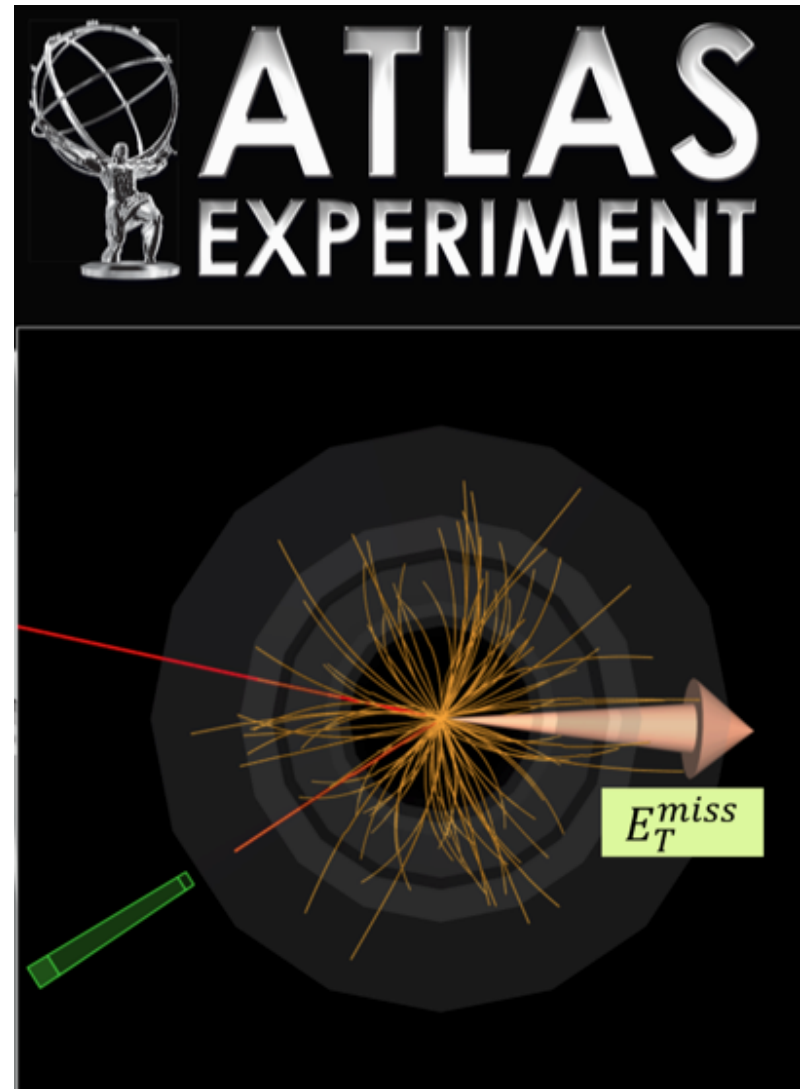


# Missing Transverse Energy

- In hadron collisions a significant unmeasured amount of energy 'escapes' in z (=beam) direction
- total momentum in transverse direction is 0, conserved in collision
- missing total  $p_T$  or  $E_T$  points to weakly or non-interacting particles
  - neutrinos, new physics

$$E_T^{\text{miss}} = -\sum_{\text{particles}}(E_T)$$

- simple strategy: sum up calo clusters and energy of escaped muons in MS
- best strategy: take calibrated physics objects, overlap removal, add unassociated clusters
  - final, 'refined MET'



Part 2, slide 41 (WW event)



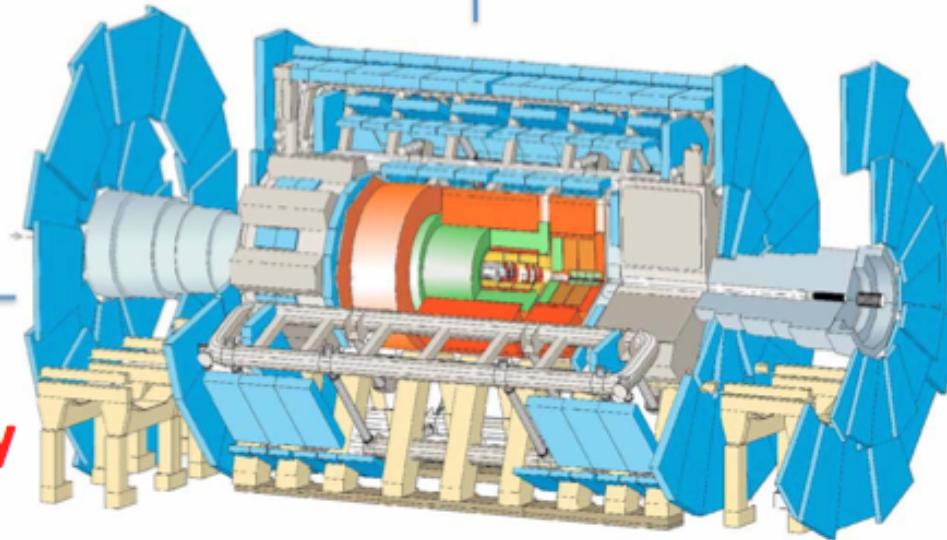
# Inputs to Refined Missing ET

## Electrons / Photons / Jets / Taus

- Overlap resolution needed for calorimeter-based signals
- Object quality cuts change MET
- Use best calibration for each

## Muons

- Use good reconstructed muons
- Possible source of fake MET
- Avoid double-counting signal in calorimeters



## Remaining Clustered Energy

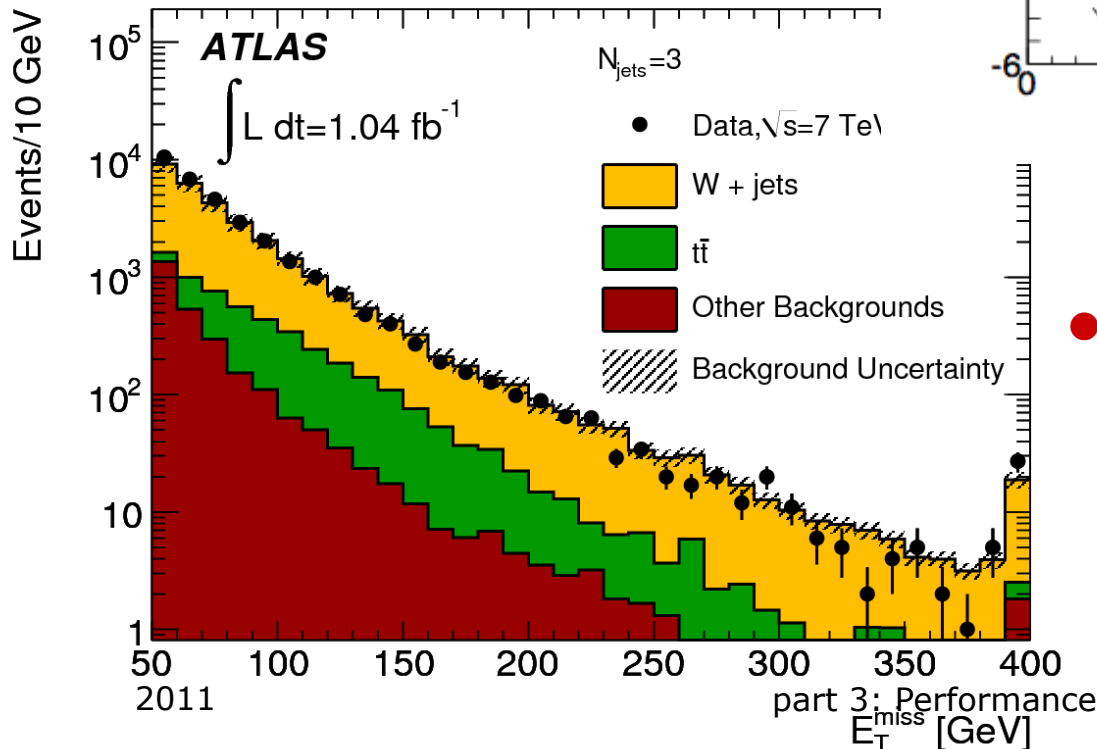
- Important to use all real signals in calorimeters, but ignore noise
- Need to derive calibration for soft signals
- Improve measurement with tracks

## Data Quality/Monitoring

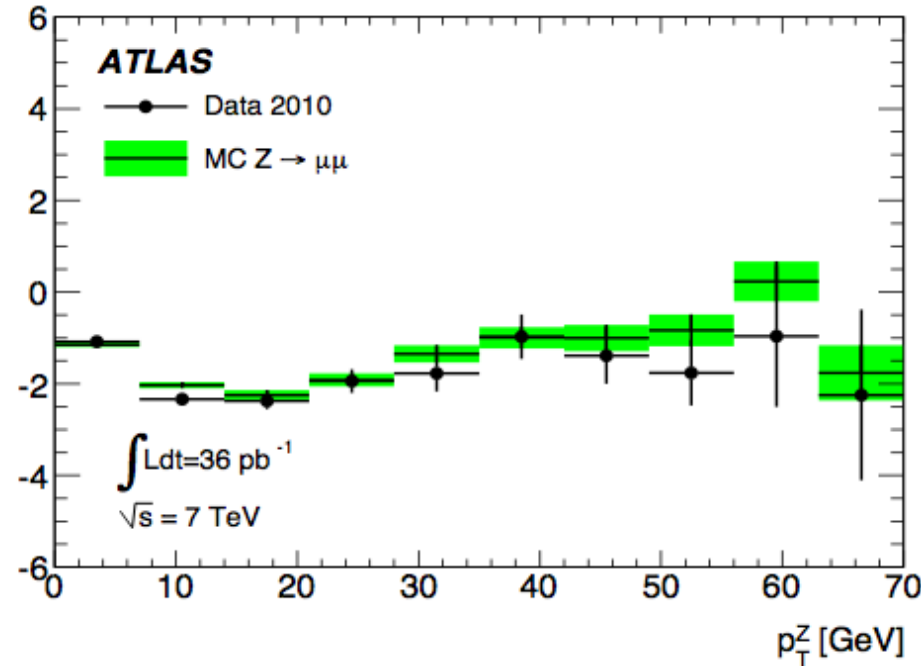
- Physics analyses must exclude/understand data with detector problems

# Missing ET Performance

- Several methods for measuring MET performance, example
  - study  $Z \rightarrow \mu\mu$  decays (MET known 0)
  - from residual bias calibrate hadronic recoil against more precisely known muon momenta



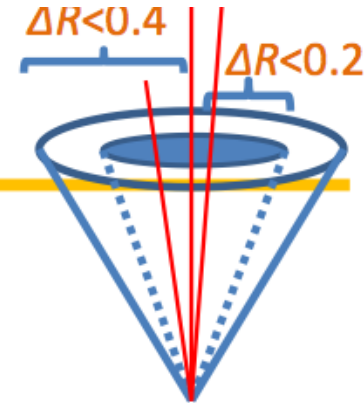
$\langle E_T^{\text{miss}} \cdot A_Z \rangle$  [GeV]



- No big surprises in inclusive MET distribution
  - detector effects vs. new physics

# Tau Reconstruction

Only hadronically decaying taus considered  
Decay to odd-numbered charged particles



- Track-seeded and **calo-seeded** candidates
  - Tracks( $p_T > 6\text{GeV}$ ) used as seed.
  - Collected tracks( $p_T > 1\text{GeV}$ ) around seed in cone  $\Delta R < 0.2$ , use them to define  $\eta, \varphi$ .
  - Look for jet (Anti-Kt algorithm with radius  $\Delta R < 0.4$  on topological clusters) around track system( $10\text{GeV}$ ,  $\Delta R < 0.2$ )
  - Collected tracks( $p_T > 1\text{GeV}$ ) around seed in cone  $\Delta R < 0.2$ .
  - Reconstruct  $\pi^0$  subclusters
  - Calorimetric  $E_T$  with H1 calibration,  $E_T^{\text{flow}}$  from tracks and calo.

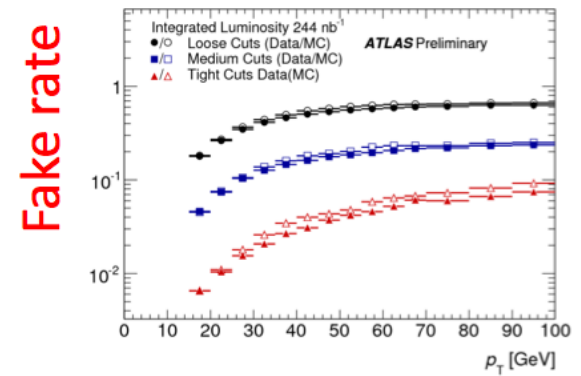
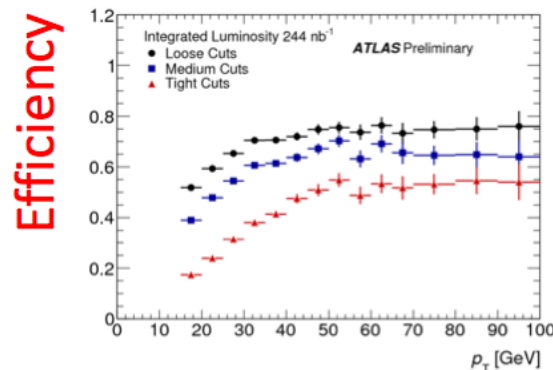


# Tau Reconstruction

- Two more categories with only one of the two seeding strategies
- Calo-seeded only candidates
  - jet seed (not yet used in Calo+Track seeded)
  - collected tracks ( $p_T > 1\text{GeV}$ ) around seed in cone  $R < 0.2$
  - calorimetric  $E_\tau$  with calibration
- Track-seeded only candidates
  - only a few % of all tau candidates
- Large number of identification variables form set of discriminators
  - including tau veto when overlap with electron/muons
- Commissioning: Only three variables are used.
  - Electro Magnetic Radius
  - Track Radius
  - $p_\tau/E_\tau$

Cut optimization (TMVA)

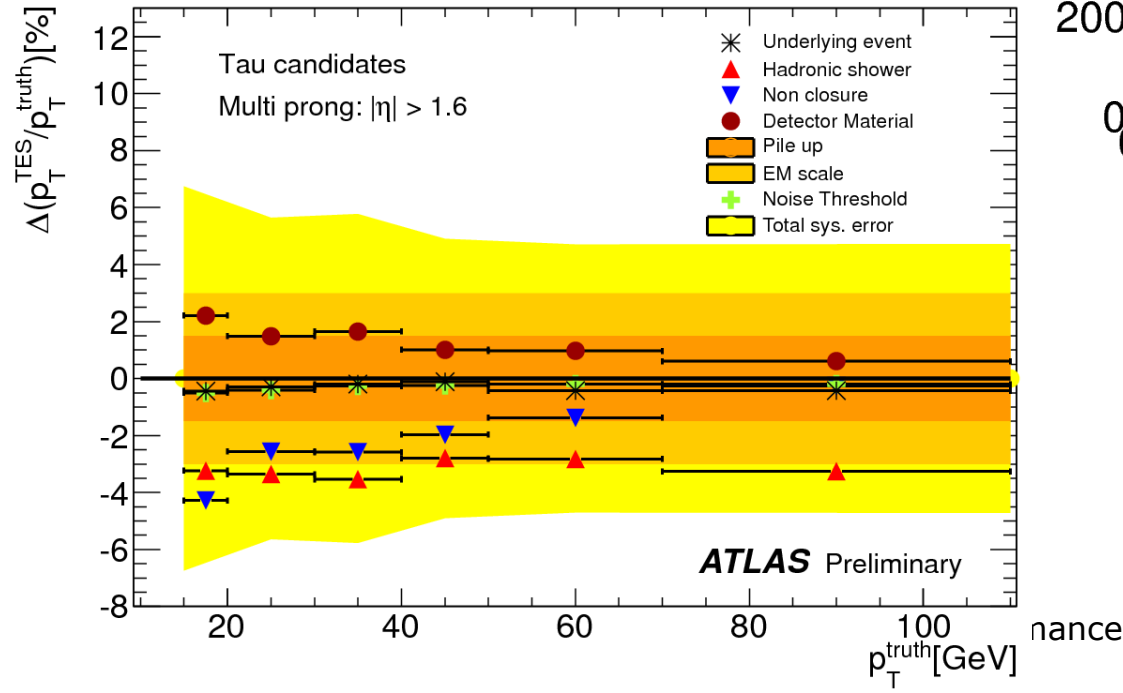
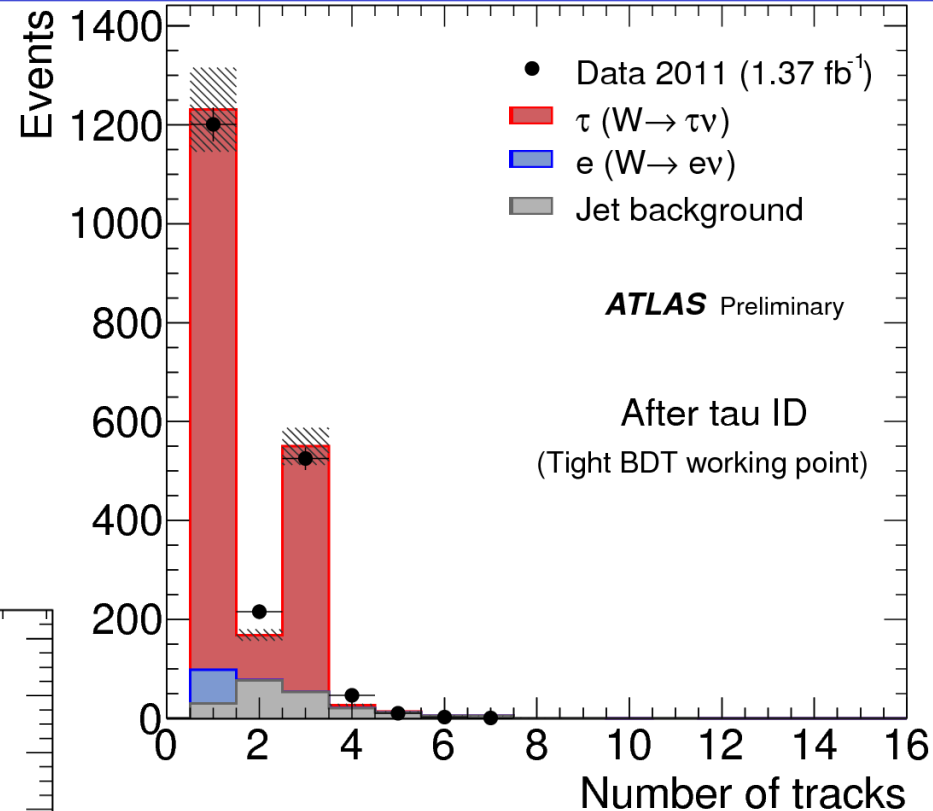
For 30%(tight), 50%(medium), 70%(loose) efficiency.



2011

# Tau Identification Performance

- Now discriminator formed from full list of variables
  - multi-variate method: boosted decision tree (BDT)
  - high purity in  $W \rightarrow \tau\nu$  decays
- first studies of  $\tau$  momentum resolution





# Summary

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- Discussed the different physics objects in ATLAS
  - outline of principal identification algorithms
  - methods to determine and improve performance on data
  - current performance, differences to simulation
- Found a remarkable performance close to design precision almost everywhere
  - only after 2 years of data-taking, previous experiments needed more time
  - a lot of effort has gone/is going into object identification and performance needs to be understood as part of physics analysis
- Last lecture will go into practical details



# Further Reading

Track and Vertex reconstruction, R. Frühwirth and A. Strandlie, Rev.Mod.Phys 82 1419 (2010)  
[http://rmp.aps.org/abstract/RMP/v82/i2/p1419\\_1](http://rmp.aps.org/abstract/RMP/v82/i2/p1419_1)

Tracking Performance Results  
<https://twiki.cern.ch/twiki/bin/view/AtlasPublic/InDetTrackingPerformanceApprovedPlots>

ATLAS conference notes  
<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES>

ATLAS papers  
<https://twiki.cern.ch/twiki/bin/view/AtlasPublic>  
– Electron Paper <http://arxiv.org/abs/1110.3174>  
– Missing ET Paper <http://arxiv.org/abs/1108.5602>

Identification of b-jets and..., N.G. Piacquadio, CERN-Thesis-2010-024  
<https://cdsweb.cern.ch/record/1243771>

The anti-kt jet clustering algorithm, M. Cacciari et al, arxiv: 0802.1189 [hep-ph]  
<http://arxiv.org/abs/0802.1189>

